# Estimates on periodic and Dirichlet eigenvalues for the Zakharov-Shabat system 

B. Grébert ${ }^{1}$, T. Kappeler ${ }^{2}$

December 19, 2002

1. UMR 6629 CNRS, Universite de Nantes, 2 rue de la Houssière, BP 92208, 44322 Nantes cedex 3, France.
2. Institut für Mathematik, Universität Zürich, Winterthurerstrasse 190, CH-8057 Zürich, Switzerland.

Abstract
Consider the $2 \times 2$ first order system due to Zakharov-Shabat,

$$
L Y:=i\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right) Y^{\prime}+\left(\begin{array}{cc}
0 & \psi_{1} \\
\psi_{2} & 0
\end{array}\right) Y=\lambda Y
$$

with $\psi_{1}, \psi_{2}$ being complex valued functions of period one in the weighted Sobolev space $H^{w} \equiv H_{\mathbb{C}}^{w}$. Denote by $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right)$ the set of periodic eigenvalues of $L\left(\psi_{1}, \psi_{2}\right)$ with respect to the interval [ 0,2 ] and by $\operatorname{spec}_{\text {Dir }}\left(\psi_{1}, \psi_{2}\right)$ the set of Dirichlet eigenvalues of $L\left(\psi_{1}, \psi_{2}\right)$ when considered on the interval $[0,1]$. It is well known that $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right)$ and $\operatorname{spec}_{\text {Dir }}\left(\psi_{1}, \psi_{2}\right)$ are discrete.

Theorem Assume that $w$ is a weight such that, for some $\delta>0$, $w_{-\delta}(k)=(1+|k|)^{-\delta} w(k)$ is a weight as well. Then for any bounded subset $\mathbb{B}$ of 1-periodic elements in $H^{w} \times H^{w}$ there exist $N \geq 1$ and $M \geq$ 1 so that for any $|k| \geq N$, and $\left(\psi_{1}, \psi_{2}\right) \in \mathbb{B}$, the set $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right) \cap\{\lambda \in$
$\mathbb{C}||\lambda-k \pi|<\pi / 2\}$ contains exactly one isolated pair of eigenvalues $\left\{\lambda_{k}^{+}, \lambda_{k}^{-}\right\}$and $\operatorname{spec}_{D i r}\left(\psi_{1}, \psi_{2}\right) \cap\left\{\lambda \in \mathbb{C}| | \lambda-k \pi \left\lvert\,<\frac{\pi}{2}\right.\right\}$ contains a single Dirichlet eigenvalue $\mu_{k}$. These eigenvalues satisfy the following estimates
(i) $\sum_{|k| \geq N} w(2 k)^{2}\left|\lambda_{k}^{+}-\lambda_{k}^{-}\right|^{2} \leq M$;
(ii) $\sum_{|k| \geq N} w(2 k)^{2}\left|\frac{\left(\lambda_{k}^{+}+\lambda_{k}^{-}\right)}{2}-\mu_{k}\right|^{2} \leq M$.

Furthermore $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right) \backslash\left\{\lambda_{k}^{ \pm},|k| \geq N\right\}$ and $\operatorname{spec}_{D i r}\left(\psi_{1}, \psi_{2}\right) \backslash\left\{\mu_{k} \mid\right.$ $|k| \geq N\}$ are contained in $\{\lambda \in \mathbb{C}||\lambda|<N \pi-\pi / 2\}$ and its cardinality is $4 N-2$, respectively $2 N-1$.
When $\psi_{2}=\bar{\psi}_{1}$ (respectively $\left.\psi_{2}=-\bar{\psi}_{1}\right), L\left(\psi_{1}, \psi_{2}\right)$ is one of the operators in the Lax pair for the defocusing (resp. focusing) nonlinear Schrödinger equation.

## Contents

1 Introduction ..... 3
1.1 Results ..... 3
1.2 Comments ..... 5
1.3 Method of proof ..... 7
1.4 Related work ..... 8
2 Periodic eigenvalues ..... 8
2.1 Lyapunov-Schmidt decomposition ..... 8
$2.2 \quad P$-equation ..... 11
2.3 $Q$-equation ..... 17
2.4 Estimates for $\alpha(n, z)$ ..... 19
2.5 Estimates for $\beta^{ \pm}(n, z)$ ..... 20
$2.6 z$-equation ..... 24
$2.7 \quad \zeta$-equation ..... 25
2.8 Proof of Theorem 1.1 ..... 26
2.9 Improvement of Theorem 1.1 for $L$ selfadjoint ..... 30
3 Riesz spaces and normal form of $L$ ..... 32
3.1 Riesz spaces ..... 32
3.2 Normal form of $L$ ..... 34
4 Dirichlet eigenvalues ..... 36
4.1 Dirichlet boundary value problem ..... 36
4.2 Decomposition ..... 38
4.3 Proof of Theorem 1.2 ..... 40
A Appendix A: Spectral properties of $L\left(\psi_{1}, \psi_{2}\right)$ ..... 41
B Appendix B: Proof of Lemma 2.8 ..... 44

## 1 Introduction

### 1.1 Results

Consider the Zakharov-Shabat operator (see [ZS])

$$
L\left(\psi_{1}, \psi_{2}\right):=i\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right) \frac{d}{d x}+\left(\begin{array}{cc}
0 & \psi_{1} \\
\psi_{2} & 0
\end{array}\right)
$$

where $\psi_{1}, \psi_{2}$ are 1-periodic elements in the weighted Sobolev space $H^{w} \equiv H_{\mathbb{C}}^{w}$ of 2-periodic functions

$$
H^{w}:=\left\{f(x)=\sum_{-\infty}^{\infty} \hat{f}(k) e^{i \pi k x} \mid\|f\|_{w}<\infty\right\}
$$

with

$$
\|f\|_{w}:=\left(2 \sum_{k \in \mathbb{Z}} w(k)^{2}|\hat{f}(k)|^{2}\right)^{1 / 2}
$$

and $w=(w(k))_{k \in \mathbb{Z}}$ a weight, i.e. a sequence of positive numbers with $w(k) \geq$ $1, w(-k)=w(k)(\forall k \in \mathbb{Z})$ and the following submultiplicative property

$$
w(k) \leq w(k-j) w(j) \quad \forall k, j \in \mathbb{Z}
$$

As an example of such a weight we mention the Sobolev weights $s_{N} \equiv$ $\left(s_{n}(k)\right)_{k \in \mathbb{Z}}, s_{N}(k):=\langle k\rangle^{N}$, where, for convenience,

$$
\langle k\rangle:=1+|k|,
$$

or more generally, the Abel-Sobolev weight $w_{a, b} \equiv\left(w_{a, b}(k)\right)_{k \in \mathbb{Z}}$

$$
w_{a, b}(k):=\langle k\rangle^{a} e^{b|k|} \quad(a \geq 0 ; b \geq 0) .
$$

An element $\psi \in H^{w_{a, b}}$ is a complex valued function $f(x)=\sum_{k \in \mathbb{Z}} \hat{f}(k) e^{i \pi k x}$, which admits an analytic extension $f(x+i y)$ to the strip $|y|<\frac{b}{\pi}$ such that $f\left(x+i \frac{b}{\pi}\right)$ and $f\left(x-i \frac{b}{\pi}\right)$ are both in the Sobolev space $H_{\mathbb{C}}^{a} \equiv H^{a}\left(\mathcal{S}^{1} ; \mathbb{C}\right)$. Denote by $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right)$ the periodic spectrum of $L\left(\psi_{1}, \psi_{2}\right)$ when considered on the interval $[0,2]$ and by $\operatorname{spec}_{\text {Dir }}\left(\psi_{1}, \psi_{2}\right)$ the Dirichlet spectrum of $L\left(\psi_{1}, \psi_{2}\right)$ when considered on $[0,1]$. It is well known that $\operatorname{both}, \operatorname{spec}\left(\psi_{1}, \psi_{2}\right)$ and $\operatorname{spec}_{D i r}\left(\psi_{1}, \psi_{2}\right)$ are discrete.
The main purpose of this paper is to study the asymptotics of the large (in absolute value) eigenvalues in $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right)$ and $\operatorname{spec}_{D i r}\left(\psi_{1}, \psi_{2}\right)$ for 1-periodic functions $\psi_{1}, \psi_{2}$ in $H^{w}$. To formulate our first result we need to introduce some more notation: we say that $w$ is a $\delta$-weight for $\delta>0$ if

$$
w_{*}(k):=\langle k\rangle^{-\delta} w(k)
$$

is a weight as well. Notice that the Abel-Sobolev weight $w_{a, b}$ is a $\delta$-weight iff $0<\delta \leq a$. Let

$$
\delta_{*}:=\delta \wedge \frac{1}{2}\left(=\inf \left(\delta, \frac{1}{2}\right)\right) .
$$

Further let

$$
\rho_{n}:=\left(\left(\hat{\psi}_{2}(2 n)+\beta_{0}^{+}(n)\right)\left(\hat{\psi}_{1}(-2 n)+\beta_{0}^{-}(n)\right)\right)^{1 / 2}
$$

with an arbitrary, but fixed choice of the square root and

$$
\begin{aligned}
& \beta_{0}^{+}(n):=\sum_{k, j \neq n} \frac{\hat{\psi}_{2}(k+n)}{(k-n) \pi} \frac{\hat{\psi}_{1}(-k-j)}{(j-n) \pi} \hat{\psi}_{2}(j+n) \\
& \beta_{0}^{-}(n):=\sum_{k, j \neq n} \frac{\hat{\psi}_{1}(-k-n)}{(k-n) \pi} \frac{\hat{\psi}_{2}(k+j)}{(j-n) \pi} \hat{\psi}_{1}(-j-n)
\end{aligned}
$$

The first result concerns the periodic eigenvalues (cf. section 2).

Theorem 1.1 Let $M \geq 1, \delta>0$ and $w$ a $\delta$-weight. Then there exist constants $1 \leq C<\infty$ and $1 \leq N<\infty$ so that the following statements hold:
For any $|n| \geq N$ and any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$, the set $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right) \cap\left\{\lambda \in \mathbb{C}\left||\lambda-n \pi|<\frac{\pi}{2}\right\}\right.$ contains exactly one isolated pair of eigenvalues $\left\{\lambda_{k}^{+}, \lambda_{k}^{-}\right\}$. These eigenvalues satisfy
(i) $\sum_{|n| \geq N} w(2 n)^{2}\left|\lambda_{n}^{+}-\lambda_{n}^{-}\right|^{2} \leq C$;
(ii) $\sum_{|n| \geq N}\langle n\rangle^{3 \delta_{*}} w(2 n)^{2} \min _{ \pm}\left|\left(\lambda_{n}^{+}-\lambda_{n}^{-}\right) \pm 2 \rho_{n}\right|^{2} \leq C$;
(iii) $\operatorname{spec}\left(\psi_{1}, \psi_{2}\right) \backslash\left\{\lambda_{n}^{ \pm}| | n \mid \geq N\right\}$ is contained in $\left\{\lambda \in \mathbb{C}\left||\lambda|<N \pi-\frac{\pi}{2}\right\}\right.$ and its cardinality is $4 N-2$.

Theorem 1.2 Let $M \geq 1, \delta>0$ and $w$ be a $\delta$-weight. Then there exist constants $1 \leq C<\infty$ and $N \leq N^{\prime}<\infty$ (with $N$ given by Theorem 1.1) so that the following statements hold:
For any $|n| \geq N^{\prime}$ and any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq$ $M$, the set $\operatorname{spec}_{D i r}\left(\psi_{1}, \psi_{2}\right) \cap\left\{\lambda \in \mathbb{C}| | \lambda-n \pi \left\lvert\,<\frac{\pi}{2}\right.\right\}$ contains exactly one eigenvalue denoted by $\mu_{n}$. These eigenvalues satisfy:
(i) $\sum_{|n| \geq N^{\prime}} w(2 n)^{2}\left|\mu_{n}-\lambda_{n}^{+}\right|^{2} \leq C$;
(ii) $\operatorname{spec}_{\text {Dir }}\left(\psi_{1}, \psi_{2}\right) \backslash\left\{\mu_{n}| | n \mid \geq N^{\prime}\right\}$ is contained in $\left\{\lambda \in \mathbb{C}\left||\lambda|<N^{\prime} \pi-\frac{\pi}{2}\right\}\right.$ and its cardinality is $2 N^{\prime}-1$.

Statement (iii) in Theorem 1.1 and (ii) in Theorem 1.2 are obtained in a standard way. For the convenience of the reader we prove it in Appendix A. In section 3, we consider the Riesz spaces $E_{n}$, i.e. the images of the Riesz projectors associated to $L\left(\psi_{1}, \psi_{2}\right)$ for a small circle around $n \pi$ with $|n|$ sufficiently large. We analyze the restriction of $L-\lambda_{n}^{+}$to $E_{n}$ and study the asymptotic properties of eigenfunctions in $E_{n}$ for $|n| \rightarrow \infty$.

### 1.2 Comments

Operator $L\left(\psi_{1}, \psi_{2}\right)$ : The Zakharov-Shabat operator occurs in the Lax pair representation $\frac{d M_{ \pm}}{d t}=\left[M_{ \pm}, A_{ \pm}\right]$of the focusing ( $N L S_{-}$) and defocusing $\left(N L S_{+}\right)$nonlinear Schrödinger equation

$$
i \partial_{t} \varphi=-\partial_{x}^{2} \pm 2|\varphi|^{2} \varphi
$$

$$
M_{+}:=L(\varphi, \bar{\varphi}) ; \quad M_{-}:=L(\varphi,-\bar{\varphi})
$$

(whereas the operators $A_{ \pm}$are rather complicated third order operators, given in [FT]). One can show that spec $L(\varphi, \bar{\varphi})$ respectively spec $L(\varphi,-\bar{\varphi})$ is a complete set of conserved quantities for $N L S_{+}$respectively $N L S_{-}$. We mention that $L\left(\psi_{1}, \psi_{2}\right)$ is unitarily equivalent to the $A K N S$ operator (see [AKNS], [MA])

$$
L_{A K N S}:=\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right) \frac{d}{d x}+\left(\begin{array}{cc}
-q & p \\
p & q
\end{array}\right)
$$

where

$$
\psi_{1}:=-q+i p ; \quad \psi_{2}=-q-i p .
$$

Hence the selfadjoint operator $M_{+}$corresponds to an operator $L_{A K N S}$ with the functions $q, p$ being real valued.

Selfadjoint case: We emphasize that Theorem 1.1 and Theorem 1.2 do not require $L\left(\psi_{1}, \psi_{2}\right)$ be selfadjoint. However, in the selfadjoint case, the decay rate of the asymptotics in Theorem 1.1 (ii) can be improved from $3 \delta_{*}$ to $4 \delta_{*}$,

$$
\sum_{|n| \geq N}\langle n\rangle^{4 \delta_{*}} w(2 n)^{2} \min _{ \pm}\left|\left(\lambda_{n}^{+}-\lambda_{n}^{-}\right) \pm 2 \rho_{n}\right|^{2} \leq M
$$

(This is proved in section 2.9).
$L^{2}$-case: Theorem 1.1 (i) and Theorem 1.2 (i) no longer hold for $H^{w}=L^{2}$ (i.e. $w(k)=1 \forall k \in \mathbb{Z}$ ) as the number $N$ in Theorem 1.1 cannot be chosen uniformly for 1-periodc functions $\psi_{1}, \psi_{2} \in L^{2}$ in a $L^{2}$-bounded set. This can be easily deduced from the examples considered by Li - McLaughlin [LM] : Assume that Theorem 1.1 (i) holds for $L^{2}$. Given $M>0$, choose $N$ as in Theorem 1.1 and $\psi_{1}, \psi_{2} \in L^{2}$ with $\left\|\psi_{j}\right\| \equiv\left\|\psi_{j}\right\|_{L^{2}}=M$. Define $\left(\psi_{1, k}, \psi_{2, k}\right)=\left(e^{2 \pi i k x} \psi_{1}, e^{-2 \pi i k x} \psi_{2}\right)(k \in \mathbb{Z})$. Then $\left\|\psi_{j, k}\right\|_{L^{2}}=\left\|\psi_{j}\right\|_{L^{2}}(\forall k)$ and, for $n \geq N, k \geq 0$

$$
\lambda_{n+k}^{ \pm}\left(\psi_{1, k}, \psi_{2, k}\right)=\lambda_{n}^{ \pm}\left(\psi_{1}, \psi_{2}\right)+k \pi
$$

which leads for appropriate choices of $\psi_{1}, \psi_{2}$ to a contradiction. For $L$ selfadjoint, a local version of Theorem 1.1 and Theorem 1.2 have been established, using different methods, in [GG]. Most likely, the analysis presented in this paper can be used to obtain a local version of Theorem 1.1 (i) and Theorem 1.2 (i) for $L$ arbitrary.

Submultiplicative property of weights: Notice that the requirement of a weight to be submultiplicative excludes weights of super-exponential growth $\exp \left(a|k|^{\alpha}\right)$ with $\alpha>1$. Most likely, the conclusions of Theorem 1.1 and Theorem 1.2 do not hold for such weights (cf. [KM] for the case of Schrödinger operators).

Boundary conditions: Similarly as in [KM] the method for proving Theorem 1.2 can be applied to a whole class of boundary conditions (cf. section 4 in $[\mathrm{KM}]$ where this class has been described for the Schrödinger operator $\left.-\frac{d^{2}}{d x^{2}}+V\right)$.
Smoothness vs. decay of gap length: For selfadjoint Zakharov-Shabat operators $L(\psi, \bar{\psi})$, Theorem 1.1 has a partial inverse. In this case, the eigenvalues $\left(\lambda_{n}^{ \pm}\right)_{n \in \mathbb{Z}}=\operatorname{spec} L(\psi, \bar{\psi})$ are real and can be ordered such that

$$
\ldots \leq \lambda_{n-1}^{+}<\lambda_{n}^{-} \leq \lambda_{n}^{+}<\lambda_{n+1}^{-} \leq \ldots ; \quad \lambda_{n}^{ \pm}=n \pi+o(1) .
$$

Given a weight $w$ and $K \geq 0$, denote by $w_{K}$ the weight $w_{K}(n):=\langle n\rangle^{K} w(n)$.

Proposition 1.3 Let $w$ be a $\delta$-weight for some $\delta>0, K \geq 0$ and $\varphi \in H^{w}$. Then $\varphi \in H^{w_{K}}$ iff

$$
\sum_{n \in \mathbb{Z}} w_{K}(2 n)^{2}\left|\lambda_{n}^{+}-\lambda_{n}^{-}\right|^{2}<\infty
$$

where $\lambda_{n}^{ \pm} \equiv \lambda_{n}^{ \pm}(\varphi, \bar{\varphi})$.

In the non selfadjoint case, the smoothness is not characterized by properties of the periodic spectrum alone (cf. [ST] for an analysis in the case of Schrödinger operators).

### 1.3 Method of proof

Typically, asymptotic estimates on the gap's lengths $\left(\lambda_{k}^{+}-\lambda_{k}^{-}\right)_{k \in \mathbb{Z}}$ of $\operatorname{spec}\left(L\left(\psi_{1}, \psi_{2}\right)\right)$ are obtained from asymptotic expansions of the eigenvalues $\lambda_{k}^{ \pm}=k \pi+\frac{c_{-1}}{k}+\ldots$ (cf. e.g. [Ma]). This approach, however does not allow to obtain the results of Theorem 1.1 and Theorem 1.2 for weights with exponential decay such as the Abel-Sobolev weight. The new feature in the proof of our results is to use as in $[\mathrm{KM}]$ a Lyapunov-Schmidt type decomposition described in detail in section 2.1.

### 1.4 Related work

Similar results as the ones presented here for the Zakharov-Shabat operator $L\left(\psi_{1}, \psi_{2}\right)$ have been obtained previously for the Schrödinger operator $-\frac{d^{2}}{d x^{2}}+V$ in $[\mathrm{KM}]$. In this paper we document that the same methods, with adjustments, can be applied to $L$. At first sight this is astonishing, as, unlike in the case of the Schrödinger operator, the distance between adjacent pairs of eigenvalues $\left(\lambda_{n}^{+}, \lambda_{n}^{-}\right)$and $\left(\lambda_{n+1}^{+}, \lambda_{n+1}^{-}\right)$does not get unbounded for $|n| \rightarrow \infty$, a fact which was used in an essential way in [KM]. We explain in section 2.1 how this problem for $L$ can be overcome.
A weaker version of Theorem 1.1 has been reported in [GKM] (cf. also [GK]). For Sobolev weights, the asymptotics of the eigenvalues $\lambda_{n}^{ \pm}$and hence of the gap length $\gamma_{n}:=\lambda_{n}^{+}-\lambda_{n}^{-}$have been obtained in the selfadjoint case by Marchenko [Ma] (cf. also [GG], [Gre], [Mis], [LS]). In the non selfadjoint case only a few results have been known so far (see [LM], [Ta1], [Ta2]).

## 2 Periodic eigenvalues

### 2.1 Lyapunov-Schmidt decomposition

Consider the Zakharov-Shabat operator

$$
L\left(\psi_{1}, \psi_{2}\right):=i\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right) \frac{d}{d x}+\left(\begin{array}{cc}
0 & \psi_{1} \\
\psi_{2} & 0
\end{array}\right)
$$

where $\psi_{1}$ and $\psi_{2}$ are in $H^{w}$. For $\psi_{1}=\psi_{2}=0$, the periodic eigenvalues are given by $\left\{\lambda_{k}^{+}, \lambda_{k}^{-} \mid k \in \mathbb{Z}\right\}$ with $\lambda_{k}^{+}=\lambda_{k}^{-}=k \pi$ and an orthonormal basis of corresponding eigenfunctions in $L^{2}[0,2] \times L^{2}[0,2]$ are given by

$$
\begin{equation*}
e_{k}^{+}(x)=\frac{1}{\sqrt{2}}\binom{0}{1} e^{i k \pi x}, e_{k}^{-}(x)=\frac{1}{\sqrt{2}}\binom{1}{0} e^{-i k \pi x} \tag{2.1}
\end{equation*}
$$

Considering the multiplication operator $\left(\begin{array}{cc}0 & \psi_{1} \\ \psi_{2} & 0\end{array}\right)$ as a perturbation of the
Dirac operator $i\left(\begin{array}{cc}1 & 0 \\ 0 & -1\end{array}\right) \frac{d}{d x}$ we will see that for $k$ sufficiently large $L$ has a
pair of eigenvalues near $k \pi$, isolated from the remaining part of the spectrum of $L$. Our aim is to obtain an estimate for the distance between the two eigenvalues and to compare the eigenvalues and corresponding eigenfunctions (or root vectors) with the corresponding ones for $\psi_{1}=\psi_{2}=0$.
We express the eigenvalue equation

$$
\begin{equation*}
L F=\lambda F \tag{2.2}
\end{equation*}
$$

in the basis $e_{k}^{+}, e_{k}^{-}(k \in \mathbb{Z})$ defined in (2.1): Given $F$ in the Sobolev space $H^{1}$, write

$$
\begin{equation*}
F(x)=\sum_{k \in \mathbb{Z}} \hat{F}_{2}(k) e_{k}^{+}(x)+\hat{F}_{1}(-k) e_{k}^{-}(x) \tag{2.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\psi_{1}(x)=\sum_{k \in \mathbb{Z}} \hat{\psi}_{1}(k) e^{i k \pi x} ; \psi_{2}(x)=\sum_{k \in \mathbb{Z}} \hat{\psi}_{2}(k) e^{i k \pi x} \tag{2.4}
\end{equation*}
$$

Substituting (2.3) - (2.4) into (2.2) leads to

$$
\begin{align*}
L F(x) & =\sum_{k \in \mathbb{Z}} k \pi\left(\hat{F}_{2}(k) e_{k}^{+}(x)+\hat{F}_{1}(-k) e_{k}^{-}(x)\right) \\
& +\sum_{k, j \in \mathbb{Z}} \hat{\psi}_{1}(-k-j) \hat{F}_{2}(j) e_{k}^{-}(x)+\hat{\psi}_{2}(k+j) \hat{F}_{1}(-j) e_{k}^{+}(x) . \tag{2.5}
\end{align*}
$$

Hence $\lambda$ is a periodic eigenvalue of $L\left(\psi_{1}, \psi_{2}\right)$, when considered on the interval $[0,2]$, iff there exists $\left(\hat{F}_{1}, \hat{F}_{2}\right) \in \ell^{2} \times \ell^{2}$ with $\left(\hat{F}_{1}, \hat{F}_{2}\right) \neq(0,0)$ such that, for all $k \in \mathbb{Z}$,

$$
\begin{align*}
& (k \pi-\lambda) \hat{F}_{2}(k)+\sum_{j \in \mathbb{Z}} \hat{\psi}_{2}(k+j) \hat{F}_{1}(-j)=0  \tag{2.6}\\
& (k \pi-\lambda) \hat{F}_{1}(-k)+\sum_{j \in \mathbb{Z}} \hat{\psi}_{1}(-k-j) \hat{F}_{2}(j)=0 \tag{2.7}
\end{align*}
$$

Here $\ell^{2} \equiv \ell^{2}(\mathbb{Z} ; \mathbb{C})$ denotes the Hilbert space of complex valued $\ell^{2}$-sequences $(a(k))_{k \in \mathbb{Z}}$. In order to solve equations (2.6) - (2.7) we consider a LyapunovSchmidt type decomposition. For $n \in \mathbb{Z}$ fixed, we look for eigenvalues near $n \pi, \lambda=n \pi+z$, with $|z| \leq \frac{\pi}{2}$. The linear system (2.6) - (2.7) is then decomposed into a two dimensional system consisting of (2.6) - (2.7) with $k=n$, referred to as the $\mathcal{Q}$-equation, and an infinite dimensional system consisting of (2.6) - (2.7) with $k \in \mathbb{Z} \backslash\{n\}$, referred to as the $\mathcal{P}$-equation.

First we introduce some more notation. For $K \in \mathbb{Z}$ and a weight $w$ denote by $\ell_{w}^{2}(K)$ the complex Hilbert space $\ell_{w}^{2}(K) \equiv \ell_{w}^{2}(K, \mathbb{C})$,

$$
\ell_{w}^{2}(K):=\left\{(a(k))_{k \in K} \mid\|a\|_{w}<\infty\right\}
$$

where $\|a\|_{w}=(a, a)_{w}^{1 / 2}$ and, for $a, b \in \ell_{w}^{2}$,

$$
(a, b)_{w}:=\sum_{k \in K} w(k)^{2} \overline{a(k)} b(k)
$$

Most frequently, we will use for $K$ the set $\mathbb{Z}$ or $\mathbb{Z} \backslash n \equiv \mathbb{Z} \backslash\{n\}$. If necessary for clarity, we write $a_{K}$ for a sequence $(a(k))_{k \in K} \in \ell_{w}^{2}(K)$.
For a linear operator $A: \ell_{w_{1}}^{2}\left(K_{1}\right) \rightarrow \ell_{w_{2}}^{2}\left(K_{2}\right)$ we denote by $A(k, j)$ its matrix elements,

$$
(A a)(k):=\sum_{j \in K_{1}} A(k, j) a(j) \quad\left(k \in K_{2}\right) .
$$

Further we introduce the shift operator $S$ and an involution operator $\mathcal{J}$

$$
\begin{aligned}
& S: \ell^{2}\left(\mathbb{Z} \rightarrow \ell^{2}(\mathbb{Z}), \quad(S a)(k):=a(k+1) \quad \forall k \in \mathbb{Z}\right. \\
& J: \ell^{2}\left(\mathbb{Z} \rightarrow \ell^{2}(\mathbb{Z}),(\mathcal{J} a)(k):=a(-k) \quad \forall k \in \mathbb{Z}\right.
\end{aligned}
$$

The restriction of $S$ to $\ell_{w}^{2}(K)$ with values in $\ell_{S^{n} w}^{2}(K)$ is again denoted by $S$ and $S^{n}:=S \circ \ldots \circ S$ denotes the $n^{\prime}$ th iterate of $S$. Notice that

$$
\left\|S^{n} a\right\|_{\ell_{S^{n} w}^{2}(K)}^{2}=\sum_{k \in K} w(k+n)^{2}|a(k+n)|^{2} \leq\|a\|_{\ell_{w}^{2}(\mathbb{Z})}^{2}
$$

For $\left(\hat{F}_{2}, \hat{F}_{1}\right) \in \ell^{2} \times \ell^{2}$, write

$$
\begin{aligned}
& \hat{F}_{2}=\left(x^{F}, \breve{F}_{2}\right), x^{F}:=\hat{F}_{2}(n) ; \quad \breve{F}_{2}:=\left(\hat{F}_{2}(k)\right)_{k \in \mathbb{Z} \backslash n} \\
& \hat{F}_{1}=\left(y^{F}, J \breve{F}_{1}\right), y^{F}:=\hat{F}_{1}(-n) ; \quad \breve{F}_{1}:=\left(\hat{F}_{1}(k)\right)_{k \in \mathbb{Z} \backslash n} .
\end{aligned}
$$

Using the above introduced notation, the equations (2.6) - (2.7) read as follows:

$$
\begin{align*}
& -z x^{F}+\hat{\psi}_{2}(2 n) y^{F}+\left\langle S^{n} \hat{\psi}_{2}, J \breve{F}_{1}\right\rangle=0  \tag{2.8}\\
& \hat{\psi}_{1}(-2 n) x^{F}-z y^{F}+\left\langle S^{n} J \hat{\psi}_{1}, \breve{F}_{2}\right\rangle=0 \tag{2.9}
\end{align*}
$$

and

$$
\begin{equation*}
\binom{y^{F}\left(S^{n} \hat{\psi}_{2}\right)_{\mathbb{Z} \backslash n}}{x^{F}\left(S^{n} J \hat{\psi}_{1}\right)_{\mathbb{Z} \backslash n}}+\left(A_{n}-z\right)\binom{\breve{F}_{2}}{J \breve{F}_{1}}=0 \tag{2.10}
\end{equation*}
$$

The equations (2.8) - (2.9) together form the $\mathcal{Q}$-equation and (2.10) is the $\mathcal{P}$-equation. The operator $A_{n}$ is given by

$$
A_{n}=\left(\begin{array}{ll}
\left((k-n) \pi \delta_{k j}\right)_{k, j \in \mathbb{Z} \backslash n} & \left(\hat{\psi}_{2}(k+j)\right)_{k, j \in \mathbb{Z} \backslash n} \\
\left(\left(J \hat{\psi}_{1}(k+j)\right)_{k, j \in \mathbb{Z} \backslash n}\right. & \left((k-n) \pi \delta_{k j}\right)_{k, j \in \mathbb{Z} \backslash n}
\end{array}\right)
$$

and $\langle\cdot, \cdot \cdot\rangle \equiv\langle\cdot, \cdot\rangle_{\mathbb{Z} \backslash n}$ is defined by (no complex conjugation)

$$
\left\langle\binom{ a_{n}}{b_{n}},\binom{c_{n}}{d_{n}}\right\rangle:=\sum_{k \in \mathbb{Z} \backslash n}\left(a_{n}(k) c_{n}(k)+b_{n}(k) d_{n}(k)\right) .
$$

For $\psi_{1}=\psi_{2}=0$ and $|z| \leq \frac{\pi}{2}$, the operator $\left(z-A_{n}\right)$ is invertible as $(k \pi-$ $(n \pi-z)) \neq 0$ for $k \neq n$. By a perturbation argument we will show that $\left(z-A_{n}\right)$ can be inverted for $|z| \leq \frac{\pi}{2}$ and $|n|$ sufficiently large which then allows to solve the $\mathcal{P}$-equation (2.10) for $\left(\breve{F}_{2}, J \breve{F}_{1}\right)$ for any $x^{F}, y^{F} \in \mathbb{C}$. This solution is substituted into (2.8) - (2.9) which leads to a homogeneous linear system of two equations for $x^{F}$ and $y^{F}$ with coefficients which depend on the parameter $z$. Hence $\lambda=n \pi+z$ is a periodic eigenvalue of $L\left(\psi_{1}, \psi_{2}\right)$ iff the corresponding determinant is equal to 0 . The nature of the latter equation allows to obtain asymptotics for the difference $\lambda_{n}^{+}-\lambda_{n}^{-}$without having to compute the asymptotics of $\lambda_{n}^{+}$and $\lambda_{n}^{-}$(cf. section 2.6-2.7).

## 2.2 $P$-equation

Let us first introduce some more notation. Denote by $\Delta_{n}$ the diagonal part of $A_{n}$

$$
\Delta_{n}:=\left(\begin{array}{cc}
D_{n} & 0 \\
0 & D_{n}
\end{array}\right) ; D_{n}:=\left((k-n) \pi \delta_{k j}\right)_{k, j \in \mathbb{Z} \backslash n}
$$

and set

$$
B_{n}:=A_{n}-\Delta_{n} .
$$

Notice that for $|z| \leq \frac{\pi}{2},\left(z-\Delta_{n}\right)^{-1}$ is invertible. Hence we may introduce

$$
T_{n} \equiv T_{n, z}:=B_{n}\left(z-\Delta_{n}\right)^{-1}=\left(\begin{array}{cc}
0 & R_{n}^{(2)}  \tag{2.11}\\
R_{n}^{(1)} & 0
\end{array}\right)
$$

where $R_{n}^{(j)} \equiv R_{n, z}^{(j)}: \ell^{2}(\mathbb{Z} \backslash n) \rightarrow \ell^{2}(\mathbb{Z} \backslash n)$ are defined by

$$
\begin{equation*}
R_{n}^{(1)}(a):=J\left(\hat{\psi}_{1} *\left(z-D_{n}\right)^{-1}\right) a ; R_{n}^{(2)}(a):=\hat{\psi}_{2} * J\left(z-D_{n}\right)^{-1} a \tag{2.12}
\end{equation*}
$$

$R_{n}^{(1)}$ and $R_{n}^{(2)}$ have the following matrix representations

$$
\begin{equation*}
R_{n}^{(1)}(k, j):=\frac{\hat{\psi}_{1}(-k-j)}{z-(j-n) \pi} ; R_{n}^{(2)}(k, j):=\frac{\hat{\psi}_{2}(k+j)}{z-(j-n) \pi}(k, j \in \mathbb{Z} \backslash n) \tag{2.13}
\end{equation*}
$$

Formally, for any $x^{F}, y^{F} \in \mathbb{C}$, the $\mathcal{P}$-equation (2.10) can be solved

$$
\binom{\breve{F}_{2}}{J \breve{F}_{1}}=\left(z-A_{n}\right)^{-1}\binom{y^{F} S^{n} \hat{\psi}_{2}}{x^{F} S^{n} J \hat{\psi}_{1}}
$$

with

$$
\begin{equation*}
\left(z-A_{n}\right)^{-1}=\left(z-\Delta_{n}\right)^{-1}\left(I d-T_{n}\right)^{-1} \tag{2.14}
\end{equation*}
$$

To justify the formal considerations above it is to show that $\left(I d-T_{n}\right)$ is invertible. Unfortunately, the norm $\left\|T_{n}\right\|$ of $T_{n}$ in $\mathcal{L}\left(\ell_{S^{n} w}^{2}\right)$ (with $\ell_{S^{n} w}^{2} \equiv$ $\left.\ell_{S^{n} w}^{2}\left(\mathbb{Z} \backslash n ; \mathbb{C}^{2}\right)\right)$ does not become small as $|n| \rightarrow \infty$. However, it turns out that, assuming an additional condition on the weight, the norm of $T_{n}^{2}$ is small for $|n| \rightarrow \infty$. The invertibility of $\left(I d-T_{n}\right)$ then follows from the identity

$$
\begin{equation*}
I d=\left(I d-T_{n}\right) \circ\left(I d+T_{n}\right)\left(I d-T_{n}^{2}\right)^{-1} . \tag{2.15}
\end{equation*}
$$

Given $\varphi \in H^{w}$, denote by $\Phi_{n}$ the operator in $\mathcal{L}\left(\ell^{2}\right)\left(\right.$ with $\left.\ell^{2} \equiv \ell^{2}(\mathbb{Z} ; \mathbb{C})\right)$ defined by $\left(n \in \mathbb{Z} ; a \in \ell^{2}(\mathbb{Z} ; \mathbb{C})\right)$

$$
\left.\left(\Phi_{n} a\right) k\right):=\sum_{j \in \mathbb{Z}} \frac{\hat{\varphi}(k+j)}{\langle n-j\rangle} a(j) \quad(\forall k \in \mathbb{Z}),
$$

where $\langle k\rangle=1+|k|$.
Recall that a weight $w$ is called a $\delta$-weight $(\delta \geq 0)$ if $w_{-\delta}(k):=\langle k\rangle^{-\delta} w(k)$ is a weight. For convenience we denote the weight $w_{-\delta}$ by $w_{*}$. The two key lemmas for proving that $\lim _{n \rightarrow \infty}\left\|T_{n}^{2}\right\|=0$ are the following ones:

Lemma 2.1 Let $w$ be a $\delta$-weight with $0 \leq \delta<\frac{1}{2}$ and $n \in \mathbb{Z}$. Then there exists $C=C(\delta)$ such that

$$
\left.\left\|\Phi_{n}\right\|_{\mathcal{L}\left(\ell_{S-n_{w}}^{2} *\right.} ; \ell_{S}^{2}{ }^{n}\right) \leq C\|\varphi\|_{w}
$$

Proof For $a \in \ell_{S^{-n} w_{*}}^{2}$ and $b \in \ell_{S^{n} w}^{2}$,

$$
\begin{aligned}
& \left|\left(b, \Phi_{n} a\right)_{S^{n} w}\right| \leq \\
& \leq \sum_{j, k} w(k+n)|b(k)| w_{*}(j-n)|a(j)| w(k+j)|\hat{\varphi}(k+j)| \cdot \\
& \frac{w(k+n)}{w_{*}(j-n) w(k+j)} \frac{4}{\langle n-j\rangle} .
\end{aligned}
$$

Using that $w$ is submultiplicative, one gets
$\frac{w(k+n)}{w_{*}(j-n) w(k+j)} \leq \frac{w(n-j)}{w_{*}(j-n)}=\langle j-n\rangle^{\delta} \leq\left(4\left|n-j+\frac{1}{2}\right|\right)^{\delta} \leq 2\left|n-j+\frac{1}{2}\right|^{\delta}$.
and hence, by the Cauchy-Schwartz inequality

$$
\begin{aligned}
& \left|\left(b, \Phi_{n} a\right)_{S^{n} w}\right| \leq \\
& \leq\|b\|_{S^{n} w}\|a\|_{S^{-n} w_{*}}\left(\sum_{k, j} \frac{4|\hat{\varphi}(k+j)|^{2} w(k+j)^{2}}{\langle n-j\rangle^{2(1-\delta)}}\right)^{1 / 2} \\
& \leq C\|b\|_{S^{n} w}\|a\|_{S^{-n} w_{*}}\|\varphi\|_{w}
\end{aligned}
$$

with $C \equiv C(\delta):=\left(\sum_{k} \frac{4}{\langle k\rangle^{2-2 \delta}}\right)^{1 / 2}<\infty$ as $\delta<\frac{1}{2}$.

Lemma 2.2 Let $\delta \geq 0$, we a $\delta$-weight and $n \in \mathbb{Z}$. Then there exists $C>0$, independent of $\delta$, such that

$$
\left\|\Phi_{n}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2} ; \ell_{S-n_{w_{*}}}^{2}\right)} \leq C \frac{\|\varphi\|_{w_{*}}}{\langle n\rangle^{\delta \wedge 1}}
$$

where as usual $\delta \wedge 1=\min (1, \delta)$.
Proof For $a \in \ell_{S^{n} w}^{2}$ and $b \in \ell_{S^{-n} w_{*}}^{2}$,

$$
\begin{aligned}
& \left|\left(b, \Phi_{n} a\right)_{S^{-n} w_{*}}\right| \leq \\
& \leq \sum_{k, j} w_{*}(k-n)|b(k)| w(j+n)|a(j)| w_{*}(k+j)|\hat{\varphi}(k+j)| \\
& \frac{w_{*}(k-n)}{w(j+n) w_{*}(k+j)} \frac{4}{\langle n-j\rangle} .
\end{aligned}
$$

As $w_{*}$ submultiplicative and symmetric,

$$
w_{*}(k-n) \leq w_{*}(k+j) w_{*}(j+n)
$$

which leads to (use definition of $w_{*}$ )

$$
\left|\left(b, \Phi_{n} a\right)_{S^{-n} w_{*}}\right| \leq\|b\|_{S^{-n} w_{*}}\|a\|_{S^{n} w}\|\hat{\varphi}\|_{w_{*}}\left(\sum_{j} \frac{4}{\langle j+n\rangle^{2 \delta}} \frac{4}{\langle j-n\rangle^{2}}\right)^{1 / 2}
$$

The claimed estimate then follows from the following elementary estimate

$$
\left(\sum_{j} \frac{1}{\langle j+n\rangle^{2 \delta}} \frac{1}{\langle j-n\rangle^{2}}\right)^{1 / 2} \leq C \frac{1}{\langle n\rangle^{\delta \wedge 1}}
$$

for some $C$, independent of $\delta$.
As an application of Lemma 2.1 and 2.2 we obtain estimates for the norms of $R_{n}^{(j)}, T_{n}$ and $T_{n}^{2}$. By definition

$$
T_{n}^{2}=\left(\begin{array}{cc}
0 & R_{n}^{(2)}  \tag{2.16}\\
R_{n}^{(1)} & 0
\end{array}\right)^{2}=\left(\begin{array}{cc}
R_{n}^{(2)} R_{n}^{(1)} & 0 \\
0 & R_{n}^{(1)} R_{n}^{(2)}
\end{array}\right)
$$

and it is useful to introduce the operators

$$
\begin{equation*}
P_{n}:=R_{n}^{(2)} R_{n}^{(1)} ; \quad Q_{n}:=R_{n}^{(1)} R_{n}^{(2)} \tag{2.17}
\end{equation*}
$$

To make notation easier we write $\ell_{S^{ \pm n} w}^{2}$ for both, $\ell_{S^{ \pm n} w}^{2}(\mathbb{Z} \backslash n ; \mathbb{C})$ and $\ell_{S^{ \pm n} w}^{2}$ $\left(\mathbb{Z} \backslash n ; \mathbb{C}^{2}\right)$.

Corollary 2.3 Let $\delta \geq 0, M \geq 1$ and $w$ be a $\delta$-weight. Then, for any 1periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M \quad(j=1,2)$, the following statements hold:
(i) If $0 \leq \delta<\frac{1}{2}$, there exists $C \equiv C(\delta)>0$ so that for $1 \leq j \leq 2, n \in \mathbb{Z}$, and $|z| \leq \frac{\pi}{2}$,

$$
\begin{aligned}
& \left\|R_{n}^{(j)}\right\|_{\mathcal{L}\left(\ell_{S-n_{*}}^{2} ; \ell_{S n_{w}}^{2}\right)} \leq C M \\
& \left\|T_{n}\right\|_{\mathcal{L}\left(\ell_{S-n_{w_{*}}}^{2} ; \ell_{S^{n} w}^{2}\right)} \leq C M
\end{aligned}
$$

(ii) If $\delta \geq 0$, there exists $C>0$ such that for $1 \leq j \leq 2, n \in \mathbb{Z}$, and $|z| \leq \frac{\pi}{2}$,

$$
\begin{aligned}
& \left\|R_{n}^{(j)}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2}, \ell_{S^{-n} w_{*}}^{2}\right)} \leq \frac{C M}{\langle n\rangle^{\delta \wedge 1}} \\
& \left\|T_{n}\right\|_{\mathcal{L}\left(\ell_{S^{n} w}^{2}, \ell_{S^{-n_{w}}}^{2}\right)} \leq \frac{C M}{\langle n\rangle^{\delta \wedge 1}} .
\end{aligned}
$$

(iii) If $0 \leq \delta<\frac{1}{2}$, then there exists $C \equiv C(\delta)$ so that for $n \in \mathbb{Z}$ and $|z| \leq \frac{\pi}{2}$,

$$
\begin{aligned}
& \left\|P_{n}\right\|_{\mathcal{L}\left(\ell_{S^{n} w}^{2}\right)} \leq \frac{C M^{2}}{\langle n\rangle^{\delta}} ; \quad\left\|Q_{n}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2}\right)} \leq \frac{C M^{2}}{\langle n\rangle^{\delta}} \\
& \left.\left\|P_{n}\right\|_{\mathcal{L}\left(\ell_{S-n_{w_{*}}}^{2}\right)} \leq \frac{C M^{2}}{\langle n\rangle^{\delta}} ; \quad\left\|Q_{n}\right\|_{\mathcal{L}\left(\ell_{S-n_{w_{*}}}^{2}\right.}\right) \leq \frac{C M^{2}}{\langle n\rangle^{\delta}} .
\end{aligned}
$$

Proof The claimed estimates for $R_{n}^{(j)}(j=1,2)$ follow from Lemma 2.1 and Lemma 2.2. As $T_{n}=\left(\begin{array}{cc}0 & R_{n}^{(2)} \\ R_{n}^{(1)} & 0\end{array}\right)$, these estimates then imply the ones for $T_{n}$. The estimates in (iii) are obtained by combining the estimates in (i) and (ii) for $R_{n}^{(j)}$.

Under the assumptions of Corollary 2.3 define, for $0<\delta<\frac{1}{2}$ and $M \geq 1$,

$$
\begin{equation*}
N_{0} \equiv N_{0}(\delta, M, w):=\max \left(1,\left(2 C M^{2}\right)^{1 / \delta}\right) \tag{2.18}
\end{equation*}
$$

with $C$ given as in Corollary 2.3 (iii).

Proposition 2.4 Let $0<\delta<\frac{1}{2}, M \geq 1$ and $w$ be a $\delta$-weight. Then, for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M,|n| \geq N_{0}$ and $|z| \leq \pi / 2$,
(i)

$$
\left\|P_{n}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2}\right)} \leq \frac{1}{2} ;\left\|Q_{n}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2}\right)} \leq \frac{1}{2}
$$

(ii) $\left(I d-P_{n}\right)$ and $\left(I d-Q_{n}\right)$ are invertible and

$$
\left\|\left(I d-P_{n}\right)^{-1}\right\|_{\mathcal{L}\left(\ell_{S} n_{w}\right)} \leq 2 ; \quad\left\|\left(I d-Q_{n}\right)^{-1}\right\|_{\mathcal{L}\left(\ell_{S} n_{w}\right)} \leq 2
$$

(iii) $I d-T_{n}^{2}$ is invertible and

$$
\left\|T_{n}^{2}\right\|_{\mathcal{L}\left(\ell_{S^{n} w}^{2}\right)} \leq \frac{1}{2} ; \quad\left\|\left(I d-T_{n}^{2}\right)^{-1}\right\|_{\mathcal{L}\left(\ell_{S^{n} w}^{2}\right)} \leq 2
$$

(iv) Statements (i) - (iii) remain true if one replaces the weight $S^{n} w$ by $S^{-n} w_{*}$.

Proof (i) By Corollary 2.3 (iii) $P_{n}$ satisfies the estimate (as $0<\delta<\frac{1}{2}$ )

$$
\left\|P_{n}\right\|_{\mathcal{L}\left(\ell_{S^{n} w}^{2}\right)} \leq \frac{C M^{2}}{\langle n\rangle^{\delta}}
$$

Hence for $|n| \geq N_{0}$

$$
\left\|P_{n}\right\|_{\mathcal{L}\left(\ell_{S}^{2} n_{w}\right)} \leq \frac{C M^{2}}{\left\langle N_{0}\right\rangle^{\delta}} \leq \frac{1}{2}
$$

Similarly, one obtains $\left\|Q_{n}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2}\right)} \leq \frac{1}{2}$.
(ii) follows immediately from (i) and (iii) follows from (i) - (ii) and the identity $T_{n}^{2}=\left(\begin{array}{cc}P_{n} & 0 \\ 0 & Q_{n}\end{array}\right)$. Finally, statements (i) - (iii) for the weight $S^{-n} w_{*}$ are proved in a similar way as for $S^{n} w$.

Summarizing the results obtained in this section, we obtain, with $\|\cdot\| \equiv$ $\|\cdot\|_{\mathcal{L}\left(\ell_{S^{n}}^{2}\left(\mathbb{Z} \backslash n ; \mathbb{C}^{2}\right)\right)}:$

Corollary 2.5 Let $0<\delta<\frac{1}{2}, M \geq 1$ and $w$ be a $\delta$-weight. Then there exists $C>0$ such that, for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$ $(j=1,2),|n| \geq N_{0}$ and $|z| \leq \pi / 2$
(i) $\left\|T_{n}\right\| \leq C$;
(ii) $\left(I d-T_{n}\right)$ is invertible in $\mathcal{L}\left(\ell_{S^{n} w}^{2}\left(\mathbb{Z} \backslash n ; \mathbb{C}^{2}\right)\right)$ and $\left\|\left(I d-T_{n}\right)^{-1}\right\| \leq C$;
(iii) $\left(z-A_{n}\right)$ is invertible in $\mathcal{L}\left(\ell_{S^{n} w}^{2}\left(\mathbb{Z} \backslash n ; \mathbb{C}^{2}\right)\right)$ and $\left\|\left(z-A_{n}\right)^{-1}\right\| \leq C$.

Proof (i) Recall that $T=\left(\begin{array}{cc}0 & R_{n}^{(2)} \\ R_{n}^{(1)} & 0\end{array}\right)$. By standard convolution estimates, there exists an absolute constant $C>0$ so that for $n \in \mathbb{Z}$ and $|z| \leq \pi / 2,\left\|\psi_{j}\right\|_{w} \leq M$

$$
\left\|T_{n}\right\| \leq C M
$$

Therefore (ii) and (iii) follow immediately from (2.14) - (2.15) and Proposition 2.4.

## 2.3 $Q$-equation

Using the notations introduced in section 2.2, we have for $|n| \geq N_{0}$ and $|z| \leq \pi / 2$

$$
\left(z-A_{n}\right)^{-1}=\left(\begin{array}{cc}
\left(z-D_{n}\right)^{-1}\left(I d-P_{n}\right)^{-1} & \left(z-D_{n}\right)^{-1} R_{n}^{(2)}\left(I d-Q_{n}\right)^{-1} \\
\left(z-D_{n}\right)^{-1} R_{n}^{(1)}\left(I d-P_{n}\right)^{-1} & \left(z-D_{n}\right)^{-1}\left(I d-Q_{n}\right)^{-1}
\end{array}\right)
$$

Hence the $P$-equation (2.10) leads to the following formulas

$$
\begin{align*}
\breve{F}_{2} & =y^{F}\left(z-D_{n}\right)^{-1}\left(I d-P_{n}\right)^{-1} S^{n} \hat{\psi}_{2}  \tag{2.19}\\
& +x^{F}\left(z-D_{n}\right)^{-1} R_{n}^{(2)}\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1} \\
J \breve{F}_{1} & =y^{F}\left(z-D_{n}\right)^{-1} R_{n}^{(1)}\left(I d-P_{n}\right)^{-1} S^{n} \hat{\psi}_{2}  \tag{2.20}\\
& +x^{F}\left(z-D_{n}\right)^{-1}\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1} .
\end{align*}
$$

These solutions are substituted into the $Q$-equation (2.8) - (2.9) to obtain for $|z| \leq \frac{\pi}{2},|n| \geq n_{0}$ the following homogeneous system

$$
\begin{gather*}
\quad\left(-z+\alpha^{+}(n, z)\right) x^{F}+\left(\hat{\psi}_{2}(2 n)+\beta^{+}(n, z)\right) y^{F}=0  \tag{2.21}\\
\left(\hat{\psi}_{1}(-2 n)+\beta^{-}(n, z)\right) x^{F}+\left(-z+\alpha^{-}(n, z)\right) y^{F}=0 \tag{2.22}
\end{gather*}
$$

where

$$
\begin{align*}
\alpha^{+}(n, z) & :=\left\langle S^{n} \hat{\psi}_{2},\left(z-D_{n}\right)^{-1}\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1}\right\rangle  \tag{2.23}\\
\beta^{+}(n, z): & =\left\langle S^{n} \hat{\psi}_{2},\left(z-D_{n}\right)^{-1} R_{n}^{(1)}\left(I d-P_{n}\right)^{-1} S^{n} \hat{\psi}_{2}\right\rangle  \tag{2.24}\\
\alpha^{-}(n, z) & :=\left\langle S^{n} J \hat{\psi}_{1},\left(z-D_{n}\right)^{-1}\left(I d-P_{n}\right)^{-1} S^{n} \hat{\psi}_{2}\right\rangle  \tag{2.25}\\
\beta^{-}(n, z): & =\left\langle S^{n} J \hat{\psi}_{1},\left(z-D_{n}\right)^{-1} R_{n}^{(2)}\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1}\right\rangle . \tag{2.26}
\end{align*}
$$

Notice that $\alpha^{ \pm}(n, z)$ and $\beta^{ \pm}(n, z)$ are analytic for $|z|<\frac{\pi}{2}$ as $R_{n}^{(j)}, P_{n}$ and $Q_{n}$ are analytic for $|z|<\frac{\pi}{2}$. An important simplification of the equations (2.21) - (2.22) results from the following observation

Lemma 2.6 For $|z| \leq \frac{\pi}{2}$ and $|n| \geq N_{0}$,

$$
\alpha^{+}(n, z)=\alpha^{-}(n, z)
$$

Proof In view of (2.23) and (2.25) it is to show that

$$
\begin{equation*}
\left(z-D_{n}\right)^{-1}\left(I d-Q_{n}\right)^{-1}=\left(\left(I d-P_{n}\right)^{-1}\right)^{t}\left(z-D_{n}\right)^{-1} \tag{2.27}
\end{equation*}
$$

where $A^{t}$ denotes the transpose of $A$,

$$
\left(A^{t}\right)(k, j):=A(j, k) \quad \text { (no complex conjugation). }
$$

The equation (2.27) can be reformulated,

$$
\left(\left(I d-Q_{n}\right)\left(z-D_{n}\right)\right)^{-1}=\left(\left(z-D_{n}\right)\left(I d-P_{n}^{t}\right)\right)^{-1}
$$

which holds iff

$$
\begin{equation*}
Q_{n}\left(z-D_{n}\right)=\left(P_{n}\left(z-D_{n}\right)\right)^{t} \tag{2.28}
\end{equation*}
$$

The identity (2.28) follows easily from

$$
\begin{aligned}
\left(Q_{n}\left(z-D_{n}\right)\right)(j, k) & =\left(R_{n}^{(1)} R_{n}^{(2)}\right)(j, k)(z-(k-n) \pi) \\
& =\sum_{\ell} \frac{\hat{\psi}_{1}(-j-\ell)}{z-(\ell-n) \pi} \cdot \hat{\psi}_{2}(\ell+k)
\end{aligned}
$$

and

$$
\begin{aligned}
\left(P_{n}\left(z-D_{n}\right)\right)^{t}(j, k) & =\left(P_{n}\left(z-D_{n}\right)\right)(k, j) \\
& =\left(R_{n}^{(2)} R_{n}^{(1)}\right)(k, j)(z-(j-n) \pi) \\
& =\sum_{\ell} \frac{\hat{\psi}_{2}(k+\ell)}{z-(\ell-n) \pi} \hat{\psi}_{1}(-\ell-j) .
\end{aligned}
$$

In view of Lemma 2.6 we write

$$
\begin{equation*}
\alpha(n, z):=\alpha^{+}(n, z) \quad\left(=\alpha^{-}(n, z)\right) . \tag{2.29}
\end{equation*}
$$

In subsequent sections we estimate the coefficients $\alpha(n, z), \beta^{+}(n, z)$ and $\beta^{-}(n, z)$.

### 2.4 Estimates for $\alpha(n, z)$

Lemma 2.7 Let $0<\delta \leq \frac{1}{2}, M \geq 1$ and $w$ be a $\delta$-weight. Then, for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M,|n| \geq N_{0}\left(N_{0} \equiv N_{0}(\delta, M)\right.$ given by (2.18)) and $|z| \leq \frac{\pi}{2}$

$$
|\alpha(n, z)| \leq \frac{4 M^{2}}{\langle n\rangle^{2 \delta}}
$$

Proof Write $\alpha(n, z)=\left\langle S^{n} \hat{\psi}_{2},\left(z-D_{n}\right)^{-1} a\right\rangle$ with $a:=\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1} \in$ $\ell_{S^{n} w}^{2}$. By Proposition 2.4,

$$
\|a\|_{S^{n} w} \leq 2\left\|\hat{\psi}_{1}\right\|_{w} \leq 2 M
$$

Hence

$$
\begin{aligned}
\langle n\rangle^{2 \delta}|\alpha(n, z)| & \leq \sum_{k \neq n} \frac{\langle n\rangle^{2 \delta}}{\langle n-k\rangle}\left|\hat{\psi}_{2}(k+n)\right||a(k)| \\
& \leq \sum_{|k+n|<|n|} \frac{\langle n\rangle^{2 \delta}}{\langle n\rangle}\left|\hat{\psi}_{2}(k+n)\right||a(k)| \\
& +\sum_{|k+n| \geq|n|} \frac{\langle n\rangle^{2 \delta}}{w(k+n)^{2}} w(k+n)\left|\hat{\psi}_{2}(k+n)\right| w(k+n)|a(k)| \\
& \leq 2\left\|\hat{\psi}_{2}\right\|_{w}\|a\|_{S^{n} w}
\end{aligned}
$$

where we used that $2 \delta \leq 1$ and $w(k+n)=w_{*}(k+n)\langle k+n\rangle^{\delta} \geq\langle n\rangle^{\delta}$ for $|k+n| \geq|n|$.

### 2.5 Estimates for $\beta^{ \pm}(n, z)$

In this section we provide estimates for $\beta^{ \pm}(n, z)$. The $\beta^{ \pm}(n, z)$ - they turn out to be quite small - determine the asymptotics of the sequence of gap lengths given in Theorem 1.1. As $\beta^{+}(n, z)(c f(2.24))$ and $\beta^{-}(n, z)(c f . ~(2.26))$ are analyzed in a similar fashion we focus on the estimate for $\beta^{+}(n, z)$. Writing $\left(I d-P_{n}\right)^{-1}=\sum_{k=0}^{\infty} P_{n}^{k}$ we obtain for $\beta^{+}(n, z)$ the following convergent series

$$
\begin{equation*}
\beta^{+}(n, z)=\sum_{k=0}^{\infty} \beta_{k}(n, z) \tag{2.30}
\end{equation*}
$$

where

$$
\begin{equation*}
\beta_{k}(n, z):=\left\langle S^{n} \hat{\psi}_{2},\left(z-D_{n}\right)^{-1} R_{n}^{(1)} P_{n}^{k} S^{n} \hat{\psi}_{2}\right\rangle \tag{2.31}
\end{equation*}
$$

The convergence of series (2.30) follows from $\left\|P_{n}\right\|_{\mathcal{L}\left(\ell_{S \pm_{n}}^{2}\right)} \leq \frac{1}{2}\left(|n| \geq N_{0}\right.$, Proposition 2.4). We begin by analyzing $\beta_{k}(n, z)$.
With $R_{n}^{(1)}$ defined by (cf (2.12))

$$
\left(R_{n}^{(1)} a\right)(j)=J\left(\hat{\psi}_{1} *\left(z-D_{n}\right)^{-1} a\right)(j)=\sum_{\ell \neq n} \frac{\left(J \hat{\psi}_{1}\right)(j+\ell) a(\ell)}{z-(\ell-n) \pi}
$$

and $\inf _{|z| \leq \frac{\pi}{2}}|z-(\ell-n) \pi| \geq \frac{1}{2}\langle\ell-n\rangle$ (for any $\left.\ell \neq n\right)$ we get

$$
\left|\left(R_{n}^{(1)} a\right)(j)\right| \leq 2 \sum_{\ell} \frac{\left|J \hat{\psi}_{1}(j+\ell)\right|}{\langle\ell-n\rangle}|a(\ell)|
$$

which leads to

$$
\begin{equation*}
\left|\beta_{k}(n, z)\right| \leq 4 \sum_{j} \frac{\left|\hat{\psi}_{2}(n+j)\right|}{\langle j-n\rangle} \sum_{\ell} \frac{\left|J \hat{\psi}_{1}(j+\ell)\right|}{\langle\ell-n\rangle}\left|\left(S^{-n} P_{n}^{k} S^{n} \hat{\psi}_{2}\right)(\ell+n)\right| \tag{2.32}
\end{equation*}
$$

Given three nonnegative sequences (i.e. sequences of nonnegative numbers), $a, b, d$ in $\ell^{2}(\mathbb{Z})$ we define, for any $n \in \mathbb{Z}$, the sequence $\Psi_{n} \equiv \Psi_{n}(a, b, d)$ by

$$
\Psi_{n}(k+n):=\sum_{j} \frac{a(k+j)}{\langle j-n\rangle} \sum_{\ell} \frac{b(j+\ell)}{\langle\ell-n\rangle} d(\ell+n)
$$

Then $\Psi_{n}$ is a nonnegative sequence in $\ell^{2}(\mathbb{Z})$ and can be used to rewrite (2.32): Introduce, for $|n| \geq N_{0}$ and $|z| \leq \frac{\pi}{2}$,

$$
\begin{equation*}
\eta_{n, 0}:=4 \Psi_{n}\left(\left|\hat{\psi}_{2}\right|,\left|J \hat{\psi}_{1}\right|,\left|\hat{\psi}_{2}\right|\right) \tag{2.33}
\end{equation*}
$$

and, for $k \geq 0$,

$$
\begin{equation*}
\eta_{n, k+1}:=4 \Psi_{n}\left(\left|\hat{\psi}_{2}\right|,\left|J \hat{\psi}_{1}\right|, \eta_{n, k}\right) \tag{2.34}
\end{equation*}
$$

where, for any $a=(a(j))_{j \in \mathbb{Z}} \in \ell^{2}(\mathbb{Z})$, we denote by $|a|$ the sequence $(|a(j)|)_{j \in \mathbb{Z}}$.

As for any $|z| \leq \frac{\pi}{2}$,

$$
\begin{aligned}
\left|\left(S^{-n} P_{n} S^{n} a\right)(k+n)\right| & =\left|\left(P_{n} S^{n} a\right)(k)\right| \\
& =\left|\left(R_{n}^{(2)} R_{n}^{(1)} S^{n} a\right)(k)\right| \\
& =\left|\left(\hat{\psi}_{2} *\left(J\left(z-D_{n}\right)^{-1}\left(R_{n}^{(1)} S^{-n} a\right)\right)\right)(k)\right| \\
& \leq 4 \sum_{j}\left|\hat{\psi}_{2}(k+j)\right| \frac{1}{\langle j-n\rangle} \sum_{\ell} \frac{\left|J \hat{\psi}_{1}(j+\ell)\right|}{\langle\ell-n\rangle}|a(\ell+n)| \\
& =4 \Psi_{n}\left(\left|\hat{\psi}_{2}\right|,\left|J \hat{\psi}_{1}\right|,|a|\right)
\end{aligned}
$$

it follows, by an induction argument, from (2.32) and

$$
S^{-n} P_{n}^{k} S^{n} \hat{\psi}_{2}=\left(S^{-n} P_{n} S^{n}\right)\left(S^{-n} P_{n}^{k-1} S^{n} \psi_{2}\right)
$$

that, for any $k \geq 0$,

$$
\begin{equation*}
\sup _{|z| \leq \frac{\pi}{2}}\left|\beta_{k}(n, z)\right| \leq \eta_{n, k}(2 n) . \tag{2.35}
\end{equation*}
$$

To estimate $\eta_{n, k}(2 n)$, we need the following auxilary lemma concerning the operator $\Psi_{n}$. For $\delta>0$ and $w$ be a $\delta$-weight, define

$$
\delta_{*}=\delta \wedge 1 / 2
$$

Lemma 2.8 Let $w$ a $\delta$-weight, and, for any $n \in \mathbb{Z}, d_{n}$ a positive sequence in $\ell_{w}^{2}$ so that

$$
\langle n\rangle^{\alpha} d_{n}(j) \leq d(j) \quad \forall n, j \in \mathbb{Z}
$$

for some $\alpha \geq 0$ and some positive sequence $d$ in $\ell_{w}^{2}$. Then there exist $C \equiv C_{\delta_{*}}$, only depending on $\delta_{*}$, and $e \in \ell_{w}^{2}$ so that for any positive sequences $a, b \in \ell_{w}^{2}$,
(i)

$$
\begin{gathered}
\sum_{n \in \mathbb{Z}}\langle n\rangle^{2\left(2 \delta_{*}+\alpha\right)} w(2 n)^{2}\left(4 \Psi_{n}\left(a, b, d_{n}\right)(2 n)\right)^{2} \\
\leq C\|a\|_{w}\|b\|_{w}\|d\|_{w} ;
\end{gathered}
$$

(ii) for any $n, j \in \mathbb{Z}$,

$$
\begin{array}{r}
\langle n\rangle^{\alpha+\delta_{*}} 4 \Psi_{n}\left(a, b, d_{n}\right)(j) \leq e(j) ; \\
\|e\|_{w} \leq C\|a\|_{w}\|b\|_{w}\|d\|_{w}
\end{array}
$$

Proof Cf. Appendix B.
From Lemma 2.8 we obtain, in view of the definition (2.33) - (2.34) and the estimate (2.35) the following

Corollary 2.9 Let $M \geq 1$ and $w$ be a $\delta$-weight. Then for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M(j=1,2)$
(i) for $k \geq 0$

$$
\sum_{|n| \geq N_{0}}\langle n\rangle^{2(2+k) \delta_{*}} w(2 n)^{2} \sup _{|z| \leq \frac{\pi}{2}}\left|\beta_{k}(n, z)\right|^{2} \leq C^{k+1} M^{2 k+3}
$$

where $1 \leq C \equiv C_{\delta}<\infty$ is given by Lemma 2.8
(ii)

$$
\sum_{|n| \geq N_{0}}\langle n\rangle^{6 \delta_{*}} w(2 n)^{2} \sup _{|z| \leq \frac{\pi}{2}}|\tilde{\beta}(n, z)|^{2} \leq C^{\prime}
$$

where $\tilde{\beta}(n, z):=\sum_{k \geq 1} \beta_{k}(n, z)$ and $1 \leq C^{\prime}<\infty$ is a constant depending only on $M$ and $\delta$.

Proof We apply Lemma 2.8 to each of the $\beta_{k}$ 's in an inductive fashion to obtain (i). Statement (ii) then follows from (i) by the Cauchy-Schwartz inequality.

To simplify further the asymptotics of $\beta$ write $\beta_{0}(n, z) \equiv \beta_{0}^{+}(n, z)=\beta_{0}^{+}(n)+$ $z \beta_{\#}^{+}(n, z)$ where

$$
\beta_{0}^{ \pm}(n):=\beta_{0}^{ \pm}(n, 0) ; \beta_{\#}^{ \pm}(n, z):=\int_{0}^{1} \partial_{z} \beta_{0}^{ \pm}(n, t z) d t
$$

As $z \mapsto \beta_{\#}^{ \pm}(n, z)$ are analytic functions in $\{|z|<\pi / 2\}$, one deduces by Cauchy's formula

$$
\begin{equation*}
\sup _{|z| \leq \pi / 4}\left|\beta_{\#}^{ \pm}(n, z)\right| \leq \frac{4}{\pi} \sup _{|z| \leq \pi / 2}\left|\beta_{0}^{ \pm}(n, z)\right| . \tag{2.36}
\end{equation*}
$$

Summarizing our results of this section gives the following

Proposition 2.10 Let $\delta>0, M \geq 1,1 \leq A<\infty$ and $w$ be a $\delta$-weight. Then there exists $C>0$ so that for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M(j=1,2)$,

$$
\begin{equation*}
\sum_{|n| \geq N_{0}}\langle n\rangle^{4 \delta_{*}} w(2 n)^{2} \sup _{|z| \leq \pi / 4}\left|\beta^{ \pm}(n, z)\right|^{2} \leq C \tag{i}
\end{equation*}
$$

(ii)

$$
\sum_{|n| \geq N_{0}}\langle n\rangle^{6 \delta_{*}} w(2 n)^{2} \sup _{|z| \leq A /\langle n\rangle_{*}}\left|\beta^{ \pm}(n, z)-\beta_{0}^{ \pm}(n)\right|^{2} \leq C .
$$

Proof Notice that $\beta^{ \pm}(n, z)=\beta_{0}^{ \pm}(n)+z \beta_{\#}^{ \pm}(n, z)+\tilde{\beta}^{ \pm}(n, z)$ and hence (i) is a consequence of Corollary 2.9 and formula (2.36). Statement (ii) is proved in the same fashion. As the supremum of $\left|\beta^{ \pm}(n, z)-\beta_{0}^{ \pm}(n)\right|$ is only taken over $|z| \leq \frac{A}{\langle n\rangle^{\delta_{*}}}$, the asymptotics of $z \beta_{\#}^{ \pm}(n, z)$ can be improved by $\delta_{*}$ to obtain from formula (2.36)

$$
\sum_{|n| \geq N_{0}}\langle n\rangle^{6 \delta_{*}} w(2 n)^{2} \sup _{|z| \leq A /\langle n\rangle^{\delta_{*}}}\left|z \beta_{\#}^{ \pm}(n, z)\right|^{2} \leq C
$$

## 2.6 z-equation

In view of (2.21) - (2.22), and (2.29), the $Q$-equation leads to the following $2 \times 2$ system

$$
\left(\begin{array}{cc}
-z+\alpha(n, z) & \hat{\psi}_{2}(2 n)+\beta^{+}(n, z)  \tag{2.37}\\
\hat{\psi}_{1}(-2 n)+\beta^{-}(n, z) & -z+\alpha(n, z)
\end{array}\right)\binom{x^{F}}{y^{F}}=\binom{0}{0}
$$

Given $|n| \geq N$ and $|z| \leq \frac{\pi}{2}$, the number $\lambda=n \pi+z$ is a periodic eigenvalue of $L$ iff there exists a nontrivial solution of $(2.37)\left(x^{F}, y^{F}\right) \in \mathbb{C}^{2} \backslash(0,0)$, or, equivalently, iff the determinant of the $2 \times 2$ matrix in (2.37) vanishes,

$$
\begin{equation*}
(z-\alpha(n, z))^{2}-\left(\hat{\psi}_{2}(2 n)+\beta^{+}(n, z)\right)\left(\hat{\psi}_{1}(-2 n)+\beta^{-}(n, z)\right)=0 . \tag{2.38}
\end{equation*}
$$

Proceeding similarly as in [KM], equation (2.38) is solved in two steps: For $\zeta$ with $|\zeta| \leq \frac{\pi}{8}$ given, consider

$$
\begin{equation*}
z_{n}=\alpha\left(n, z_{n}\right)+\zeta \tag{2.39}
\end{equation*}
$$

Substituting a solution $z(\zeta) \equiv z_{n}(\zeta)$ of (2.39) into (2.38) leads to an equation for $\zeta \equiv \zeta_{n}$,

$$
\begin{equation*}
\zeta^{2}-\left(\hat{\psi}_{2}(2 n)+\beta^{+}(n, z(\zeta))\left(\hat{\psi}_{1}(-2 n)+\beta^{-}(n, z(\zeta))\right)=0\right. \tag{2.40}
\end{equation*}
$$

Equation (2.39) is referred to as the $z$-equation and equation (2.40) as the $\zeta$-equation .

In this section we deal with the $z$-equation (2.39). To solve it we use the contractive mapping principle. According to Lemma 2.7 we can choose $N_{1} \geq$ $N_{0}$ (with $N_{0}$ given by (2.18)) so that for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$ and $|n| \geq N_{1}$

$$
\begin{equation*}
\sup _{|z| \leq \frac{\pi}{2}}|\alpha(n, z)|<\pi / 8 \tag{2.41}
\end{equation*}
$$

The following result can be proved by the same line of arguments used in the proof of [KM, Proposition 1.6].

Proposition 2.11 Let $M \geq 1,0<\delta \leq 1 / 2$ and $w$ be a $\delta$-weight. Then, there exists $N_{1} \geq N_{0}$ so that for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M,|\zeta| \leq \frac{\pi}{8}$ and $|n| \geq N_{1}$, equation (2.39) has a unique solution $z_{n}=z_{n}(\zeta)$ satisfying $\left|z_{n}\right|<\pi / 4$. The solution depends analytically on $\zeta$.

## $2.7 \quad \zeta$-equation

In this section, we improve the existence of solutions of the $\zeta$-equation (2.40)

$$
\zeta^{2}-\left(\hat{\psi}_{2}(2 n)+\beta^{+}(n, z(\zeta))\left(\hat{\psi}_{1}(-2 n)+\beta^{-}(n, z(\zeta))\right)=0\right.
$$

using Rouché's Theorem. Introduce

$$
\begin{equation*}
r_{n}:=\left(\left|\hat{\psi}_{2}(2 n)\right|+\sup _{|z| \leq \frac{\pi}{2}}\left|\beta^{+}(n, z)\right|\right) \vee\left(\left|\hat{\psi}_{1}(-2 n)\right|+\sup _{|z| \leq \frac{\pi}{2}}\left|\beta^{-}(n, z)\right|\right) . \tag{2.42}
\end{equation*}
$$

Using the same line of arguments used in the proof of [KM, Proposition 1.15] one obtains the following

Proposition 2.12 Let $M \geq 1,0<\delta \leq \frac{1}{2}$ and $w$ be a $\delta$-weight. Then there exists $N_{2} \geq N_{1}$ so that, for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$ and $|n| \geq N_{2}$, equation (2.40) has exactly two (counted with multiplicity) solutions $\zeta_{n}^{+}$, $\zeta_{n}^{-}$in $\overline{\mathcal{D}}_{r_{n}}$.

### 2.8 Proof of Theorem 1.1

In this section, Theorem 1.1 is proved.
Proof of Theorem 1.1 (i) Let $z_{n}^{ \pm}:=z\left(\zeta_{n}^{ \pm}\right)=\zeta_{n}^{ \pm}+\alpha\left(n, z_{n}^{ \pm}\right)$where $\zeta_{n}^{ \pm}$are the two solutions of the $\zeta$-equation provided by Proposition $2.12\left(|n| \geq N_{2}\right)$. Then, for $|n| \geq N_{2}$

$$
\begin{equation*}
\left|z_{n}^{+}-z_{n}^{-}\right| \leq\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right|+\sup _{|z| \leq \frac{\pi}{4}}\left|\frac{d}{d z} \alpha(n, z)\right|\left|z_{n}^{+}-z_{n}^{-}\right| . \tag{2.43}
\end{equation*}
$$

As $N_{2} \geq N_{1}$ and $|n| \geq N_{2}$ one has by the analyticity of $z \mapsto \alpha(n, z)$ and (2.41)

$$
\sup _{|z| \leq \frac{\pi}{4}}\left|\frac{d}{d z} \alpha(n, z)\right| \leq \frac{1}{2} .
$$

Together with $\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right| \leq\left|\zeta_{n}^{+}\right|+\left|\zeta_{n}^{-}\right| \leq 2 r_{n}$ equation (2.43) then leads to

$$
\left|z_{n}^{+}-z_{n}^{-}\right| \leq 4 r_{n}
$$

By the definition (2.42) of $r_{n}$, the estimates of $\beta_{n}^{ \pm}$in Proposition 2.10 (i) and the identity $\lambda_{n}^{+}-\lambda_{n}^{-}=z_{n}^{+}-z_{n}^{-}$, the latter equation implies that there exists $C \geq 1$ such that, for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w},\left\|\psi_{j}\right\|_{w} \leq M$,

$$
\sum_{|n| \geq N_{2}} w(2 n)^{2}\left|\lambda_{n}^{+}-\lambda_{n}^{-}\right|^{2} \leq C
$$

Towards the proof of Theorem 1.1 (ii), rewrite equation (2.40),

$$
\begin{equation*}
\left(\zeta_{n}^{ \pm}\right)^{2}-\rho_{n}^{2}=\eta\left(n, z\left(\zeta_{n}^{ \pm}\right)\right) \tag{2.44}
\end{equation*}
$$

where

$$
\rho_{n}=\left(\left(\hat{\psi}_{2}(2 n)+\beta_{0}^{+}(n)\right)\left(\hat{\psi}_{1}(-2 n)+\beta_{0}^{-}(n)\right)\right)^{1 / 2}
$$

with an arbitrary but fixed choice of the square root and

$$
\begin{align*}
\eta(n, z) & =\hat{\psi}_{2}(2 n)\left(\beta^{-}(n, z)-\beta_{0}^{-}(n)\right) \\
& +\hat{\psi}_{1}(-2 n)\left(\beta^{+}(n, z)-\beta_{0}^{+}(n)\right)  \tag{2.45}\\
& +\left(\beta^{-}(n, z)-\beta_{0}^{-}(n)\right)\left(\beta^{+}(n, z)-\beta_{0}^{+}(n)\right) .
\end{align*}
$$

In view of the definition (2.42) and as $w$ is assumed to be a $\delta$-weight, we have for some constant $C_{1} \geq 1$ depending on $\delta$ and $M$

$$
\begin{equation*}
r_{n} \leq \frac{C_{1}}{\langle n\rangle^{\delta_{*}}} \quad\left(\forall|n| \geq N_{0}\right) \tag{2.46}
\end{equation*}
$$

By Lemma 2.7 there exists $C_{2} \geq 1$ depending on $\delta$ and $M$ such that for $|n| \geq N_{0}$ and $|z| \leq \pi / 2$

$$
\begin{equation*}
|\alpha(n, z)| \leq \frac{C_{2}}{\langle n\rangle^{\delta_{*}}} . \tag{2.47}
\end{equation*}
$$

Let $A=C_{1}+C_{2}$ and define

$$
\begin{equation*}
s_{n}:=\sup _{|z| \leq 2 A /\langle n\rangle_{*}}|\eta(n, z)| . \tag{2.48}
\end{equation*}
$$

Notice that by Proposition 2.10 (ii), there exists $C>0$ so that

$$
\begin{equation*}
\sum_{|n| \geq N_{2}}\langle n\rangle^{3 \delta_{*}} w(2 n)^{2} s_{n} \leq C \tag{2.49}
\end{equation*}
$$

Choose $N_{3} \geq N_{2}$, depending on $\delta$ and $M$, so that

$$
\begin{equation*}
\frac{\langle n\rangle^{\delta_{*}}}{A} \sqrt{s_{n}}<\frac{1}{2}, \quad \forall|n| \geq N_{3} . \tag{2.50}
\end{equation*}
$$

Lemma 2.13 Let $M \geq 1,0<\delta$ and $w$ be a $\delta$-weight. For 1-periodic functions $\psi_{1}, \psi_{2}$ in $H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M(j=1,2)$ and $|n| \geq N_{3}$,

$$
\left|\zeta_{n}^{+}-\rho_{n}\right|+\left|\zeta_{n}^{-}+\rho_{n}\right| \leq 6 \sqrt{s_{n}}
$$

or

$$
\left|\zeta_{n}^{+}-\rho_{n}\right|+\left|\zeta_{n}^{-}-\rho_{n}\right| \leq 6 \sqrt{s_{n}} .
$$

Proof W.l.o.g. assume that $\delta \leq 1 / 2$ and hence $\delta=\delta_{*}$. By (2.44) we have for $|n| \geq N_{3}$

$$
\begin{equation*}
\left(\zeta_{n}^{ \pm}-\rho_{n}\right)\left(\zeta_{n}^{ \pm}+\rho_{n}\right)=\eta\left(n, z\left(\zeta_{n}^{ \pm}\right)\right) \tag{2.51}
\end{equation*}
$$

By definition, $z\left(\zeta_{n}^{ \pm}\right)=\zeta_{n}^{ \pm}+\alpha\left(n, z\left(\zeta_{n}^{ \pm}\right)\right)$and therefore from (2.46), (2.47) and Proposition 2.12, we conclude for $|n| \geq N_{3}$,

$$
\begin{equation*}
\left|z\left(\zeta_{n}^{ \pm}\right)\right| \leq \frac{A}{\langle n\rangle^{\delta}} \tag{2.52}
\end{equation*}
$$

From the definition of $s_{n}$ (see (2.48)) and (2.51) we deduce

$$
\begin{equation*}
\left|\zeta_{n}^{ \pm}-\rho_{n}\right|\left|\zeta_{n}^{ \pm}+\rho_{n}\right| \leq s_{n} \tag{2.53}
\end{equation*}
$$

Thus $\min _{ \pm}\left|\zeta_{n}^{+} \pm \rho_{n}\right| \leq \sqrt{s_{n}}$ and $\min _{ \pm}\left|\zeta_{n}^{-} \pm \rho_{n}\right| \leq s_{n}^{1 / 2}$. We distinguish two cases:
case $1\left|\rho_{n}\right| \leq 2 \sqrt{s_{n}}$. In this case $\left|\zeta_{n}^{ \pm}-\rho_{n}\right| \leq \sqrt{s_{n}}$ implies

$$
\left|\zeta_{n}^{ \pm}+\rho_{n}\right| \leq\left|\zeta_{n}^{ \pm}-\rho_{n}\right|+2\left|\rho_{n}\right| \leq 5 \sqrt{s_{n}}
$$

and, similarly, $\left|\zeta_{n}^{ \pm}+\rho_{n}\right| \leq \sqrt{s_{n}}$ implies $\left|\zeta_{n}^{ \pm}-\rho_{n}\right| \leq 5 \sqrt{s_{n}}$, thus Lemma 2.13 is proves in case 1 .
case $2\left|\rho_{n}\right|>2 \sqrt{s_{n}}$. It suffices to show that it is impossible to have $\max _{ \pm}\left|\zeta_{n}^{ \pm}-\rho_{n}\right| \leq \sqrt{s_{n}}$, or $\max _{ \pm}\left|\zeta_{n}^{ \pm}+\rho_{n}\right| \leq \sqrt{s_{n}}$. To the contrary, assume that

$$
\begin{equation*}
\max _{ \pm}\left|\zeta_{n}^{ \pm}-\rho_{n}\right| \leq \sqrt{s_{n}} \tag{2.54}
\end{equation*}
$$

(The other case is treated in the same way.) By (2.54), $\left|\zeta_{n}^{ \pm}+\rho_{n}\right| \geq 2\left|\rho_{n}\right|-$ $\sqrt{s_{n}}>\frac{3}{2}\left|\rho_{n}\right|$, hence

$$
\begin{equation*}
\left|\zeta_{n}^{+}+\zeta_{n}^{-}\right| \geq\left|\zeta_{n}^{+}+\rho_{n}\right|-\left|\zeta_{n}^{-}-\rho_{n}\right|>\left|\rho_{n}\right| \tag{2.55}
\end{equation*}
$$

Divide

$$
\left(\zeta_{n}^{+}\right)^{2}-\left(\zeta_{n}^{-}\right)^{2}=\eta\left(n, z\left(\zeta_{n}^{+}\right)\right)-\eta\left(n, z\left(\zeta_{n}^{-}\right)\right)
$$

by $\zeta_{n}^{+}+\zeta_{n}^{-}$and use (2.55) and (2.52) to deduce

$$
\begin{equation*}
\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right| \leq \frac{1}{\left|\rho_{n}\right|} \sup _{|z| \leq A /\langle n\rangle^{\delta}}\left|\frac{d \eta}{d z}(n, z)\right|\left|z\left(\zeta_{n}^{+}\right)-z\left(\zeta_{n}^{-}\right)\right| . \tag{2.56}
\end{equation*}
$$

To arrive at a contradiction we first show that $\zeta_{n}^{+}-\zeta_{n}^{-}=0$. As $\mid z\left(\zeta_{n}^{+}\right)-$ $z\left(\zeta_{n}^{-}\right)\left|\leq\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right|+\sup _{|z| \leq \pi / 2}\right| \frac{d}{d z} \alpha(n, z)\left|\left|z\left(\zeta_{n}^{+}\right)-z\left(\zeta_{n}^{-}\right)\right|\right.$, (2.41) leads to $\left(|n| \geq N_{2}\right)$

$$
\begin{equation*}
\left|z\left(\zeta_{n}^{+}\right)-z\left(\zeta_{n}^{-}\right)\right| \leq 2\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right| \tag{2.57}
\end{equation*}
$$

On the other hand, as $z \mapsto \eta(n, z)$ is analytic in $\{z,|z|<\pi / 2\}$, we have by Cauchy's inequality,

$$
\begin{align*}
\sup _{|z| \leq \frac{A}{\langle n\rangle^{\delta}}}\left|\frac{d}{d z} \eta(n, z)\right| & \leq \frac{\langle n\rangle^{\delta}}{A} \sup _{|z| \leq \frac{2 A}{\langle n\rangle^{\delta}}}|\eta(n, z)|  \tag{2.58}\\
& \leq \frac{\langle n\rangle^{\delta}}{A} s_{n}
\end{align*}
$$

Combining (2.56) - (2.57) with (2.50) we obtain,

$$
\begin{aligned}
\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right| & \leq \frac{2}{\left|\rho_{n}\right|} \frac{\langle n\rangle^{\delta}}{A} s_{n}\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right| \\
& \leq \frac{\langle n\rangle^{\delta}}{A} \sqrt{s_{n}}\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right| \leq \frac{1}{2}\left|\zeta_{n}^{+}-\zeta_{n}^{-}\right|
\end{aligned}
$$

and we conclude that $\zeta_{n}^{+}=\zeta_{n}^{-} \equiv \zeta_{n}$. This contradicts the assumption $\left|\rho_{n}\right|>2 \sqrt{s_{n}}$ as one can see in the following way: By the equation (2.44), $2 \zeta_{n}=\frac{d}{d \zeta} \eta\left(n, z\left(\zeta_{n}\right)\right)=\frac{d}{d z} \eta\left(n, z\left(\zeta_{n}\right)\right) \cdot \frac{d}{d \zeta} z\left(\zeta_{n}\right)$. By (2.58), $\left|\frac{d}{d z} \eta\left(n, z\left(\zeta_{n}\right)\right)\right| \leq$ $\frac{\langle n\rangle^{\delta}}{A} s_{n}$ and by $(2.41),\left|\frac{d}{d \zeta} z(\zeta)\right|=\left|\frac{d}{d \zeta}(\zeta+\alpha(n, z(\zeta)))\right| \leq 1+\frac{1}{2} \leq 2$, hence

$$
\begin{equation*}
\left|\zeta_{n}\right| \leq \frac{\langle n\rangle^{\delta}}{A} s_{n} \tag{2.59}
\end{equation*}
$$

and, by (2.55),

$$
\left|\rho_{n}\right|<2\left|\zeta_{n}\right| \leq 2 \frac{\langle n\rangle^{\delta}}{A} s_{n}<\sqrt{s_{n}} .
$$

where for the last inequality we used (2.50).
Proof of Theorem 1.1 (ii): Let $N_{3}$ be given by (2.50). Recall that

$$
\lambda_{n}^{+}-\lambda_{n}^{-}=z_{n}^{+}-z_{n}^{-}=\zeta_{n}^{+}-\zeta_{n}^{-}+\alpha\left(n, z\left(\zeta_{n}^{+}\right)\right)-\alpha\left(n, z\left(\zeta_{n}^{-}\right)\right) .
$$

By Lemma 2.13, for $|n| \geq N_{3}$,

$$
\min _{ \pm}\left|\left(\zeta_{n}^{+}-\zeta_{n}^{-}\right) \pm 2 \rho_{n}\right| \leq 6 \sqrt{s_{n}}
$$

By the analyticity of $\alpha(n, z)$ and Lemma 2.7, for $|n| \geq N_{0}$,

$$
\sup _{|z| \leq \pi / 4}\left|\frac{d}{d z} \alpha(n, z)\right| \leq \frac{C}{\langle n\rangle^{2 \delta}}
$$

Combining these two estimates, we get for $|n| \geq N_{3}$,

$$
\begin{aligned}
\min _{ \pm} \mid\left(\lambda_{n}^{+}-\lambda_{n}^{-}\right) & \pm 2 \rho_{n}\left|\leq \min _{ \pm}\right|\left(\zeta_{n}^{+}-\zeta_{n}^{-}\right) \pm 2 \rho_{n} \mid \\
& +\left(\sup _{|z| \leq \pi / 2}\left|\frac{d}{d z} \alpha(n, z)\right|\right)\left|\lambda_{n}^{+}-\lambda_{n}^{-}\right| \\
& \leq 6 \sqrt{s_{n}}+C \frac{\left|\lambda_{n}^{+}-\lambda_{n}^{-}\right|}{\langle n\rangle^{2 \delta}}
\end{aligned}
$$

Hence, by (2.49) and Theorem 1.1 (i),

$$
\sum_{|n| \geq N_{3}}\langle n\rangle^{3 \delta} w(2 n)^{2} \min _{ \pm}\left|\left(\lambda_{n}^{+}-\lambda_{n}^{-}\right) \pm 2 \rho_{n}\right|^{2} \leq C
$$

### 2.9 Improvement of Theorem 1.1 for $L$ selfadjoint

For $\psi$ a 1-periodic functions in $H^{w}$, the operator $L(\psi, \bar{\psi})$ is selfadjoint. In this section we show that in this case the decay rate of the asymptotics in Theorem 1.1 (ii) can be improved as follows :

Theorem 2.14 Let $M \geq 1, \delta>0$ and $w$ be a $\delta$-weight. Then there exist constants $1 \leq C<\infty$ and $1 \leq N<\infty$ so that for any $|n| \geq N$ and any 1-periodic function $\psi \in H^{w}$ with $\|\psi\|_{w} \leq M$,

$$
\sum_{|n| \geq N}\langle n\rangle^{4 \delta_{*}} w(2 n)^{2} \min _{ \pm}\left|\left(\lambda_{n}^{+}-\lambda_{n}^{-}\right) \pm 2 \rho_{n}\right|^{2} \leq C
$$

Proof Using the definition (2.13) with $\psi_{1}=\psi$ and $\psi_{2}=\bar{\psi}$ we get $\overline{R_{n}^{(1)}(k, j)}(\bar{z})=$ $R_{n}^{(2)}(k, j)(z)$ and thus $\overline{\beta^{-}(n, z)}=\beta^{+}(n, \bar{z})$. As the eigenvalues $\lambda_{n}^{ \pm}=n \pi+$
$z\left(\zeta_{n}^{ \pm}\right)$of $L(\psi, \bar{\psi})$ are real, equation (2.40) then reads (with $|n| \geq N_{2}$ and $N_{2}$ as in Proposition 2.12)

$$
\begin{align*}
\left(\zeta_{n}^{ \pm}\right)^{2} & =\left|\hat{\psi}(2 n)+\beta^{+}\left(n, z\left(\zeta_{n}^{ \pm}\right)\right)\right|^{2}  \tag{2.60}\\
& =\left|\hat{\psi}(2 n)+\beta_{0}^{+}(n)+\tilde{\beta}^{+}\left(n, z\left(\zeta_{n}^{ \pm}\right)\right)\right|^{2}
\end{align*}
$$

and $\rho_{n}$ is given by (with an appropriate choice of the square root)

$$
\rho_{n}=\left|\hat{\psi}(2 n)+\beta_{0}^{+}(n)\right| .
$$

Let $t_{n}:=\sup _{|z| \leq A /\langle n\rangle_{*}{ }^{\delta_{2}}}\left|\tilde{\beta}^{+}(n, z)\right|$, where $A:=C_{1}+C_{2}$ and $C_{1}$ are $C_{2}$ are defined by (2.46), (2.47). By Proposition 2.10 (ii),

$$
\begin{equation*}
\sum_{|n| \geq N_{0}}\langle n\rangle^{6 \delta_{*}} w(2 n)^{2} t_{n}^{2} \leq C \tag{2.61}
\end{equation*}
$$

From (2.60) we deduce $\min _{ \pm}\left|\zeta_{n}^{+} \pm \rho_{n}\right| \leq t_{n}$ and $\min _{ \pm}\left|\zeta_{n}^{-} \pm \rho_{n}\right| \leq t_{n}$. Substituting Lemma 2.15 below for Lemma 2.13, Theorem 2.14 follows in the same way as Theorem 1.1 (ii).

Define $N_{4} \geq N_{2}$ such that

$$
12\langle n\rangle^{\delta} t_{n}<A \quad \forall|n| \geq N_{4}
$$

Lemma 2.15 Let $M \geq 1, \delta>0$ and $w$ be a $\delta$-weight. For any 1-periodic function $\psi$ in $H^{w}$ with $\|\psi\|_{w} \leq M$ and $|n| \geq N_{4}$,

$$
\left|\zeta_{n}^{+}-\rho_{n}\right|+\left|\zeta_{n}^{-}+\rho_{n}\right| \leq 6 t_{n}
$$

or

$$
\left|\zeta_{n}^{+}+\rho_{n}\right|+\left|\zeta_{n}^{-}-\rho_{n}\right| \leq 6 t_{n}
$$

Proof The proof is similar to the one of Lemma 2.13.

## 3 Riesz spaces and normal form of $L$

### 3.1 Riesz spaces

Let $M \geq 1, \delta>0$ and $w$ be a $\delta$-weight. By Theorem 1.1, there exists $1 \leq$ $N<\infty$ so that for any 1-periodic functions $\psi_{1}, \psi_{2}$ in $H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$, the operator $L=L\left(\psi_{1}, \psi_{2}\right)$ has two (counted with multiplicity) periodic eigenvalues $\lambda_{n}^{+}, \lambda_{n}^{-}$near $n \pi$.
In Appendix A we introduce the periodic and antiperiodic boundary conditions $b c \mathrm{Per}^{+}$and $b c \mathrm{Per}^{-}$. We point out that

$$
\operatorname{spec} L=\operatorname{spec} L_{P e r}+\cup \operatorname{spec} L_{P e r^{-}}
$$

and introduce the Riesz projectors $\Pi_{2 n}: L^{2}\left([0,1] ; \mathbb{C}^{2}\right) \rightarrow L^{2}\left([0,1] ; \mathbb{C}^{2}\right)$, corresponding to $b c$ Per $^{+}$and $\Pi_{2 n-1}: L^{2}\left([0,1] ; \mathbb{C}^{2}\right) \rightarrow L^{2}\left([0,1] ; \mathbb{C}^{2}\right)$, corresponding to $b c P e r^{-}(n \in \mathbb{Z})$. Further denote by $E_{n}$ the $\mathbb{C}$-vector spaces

$$
E_{n}:=\Pi_{n}\left(L^{2}\left([0,1] ; \mathbb{C}^{2}\right)\right) \quad(|n| \geq N)
$$

Notice that $\operatorname{dim}_{\mathbb{C}} E_{n}=\operatorname{tr} \Pi_{n}=2 \forall|n| \geq N$. If $\lambda_{n}^{+} \neq \lambda_{n}^{-}$or $\lambda_{n}^{+}=\lambda_{n}^{-}$is of geometric multiplicity two, there exists a basis of $E_{n}$ consisting of eigenfunctions $F^{+}$and $F^{-}$corresponding to the eigenvalues $\lambda_{n}^{ \pm}$. If $\lambda_{n}^{+}=\lambda_{n}^{-}$is of geometric multiplicity $1, E_{n}$ is the root space of $\lambda_{n}^{+}$. Denote by $F$ a $L^{2}$-normalized eigenfunction of $L$ corresponding to the eigenvalue $\lambda=n \pi+z$,

$$
(L-\lambda) F=0, \quad\|F\|=1
$$

where $\|\cdot\|$ denotes the $L^{2}$-norm in $L^{2}\left([0,1] ; \mathbb{C}^{2}\right)$. Then

$$
F(x)=x^{F} e_{n}^{+}(x)+y^{F} e_{n}^{-}(x)+\sum_{k \neq n}\left(\breve{F}_{2}(k) e_{k}^{+}(x)+\breve{F}_{1}(-k) e_{k}^{-}(x)\right)
$$

where

$$
\binom{\breve{F}_{2}}{J \breve{F}_{1}}=\left(\begin{array}{ll}
V_{11} & V_{12} \\
V_{21} & V_{22}
\end{array}\right)\binom{y^{F}}{x^{F}}
$$

with

$$
\begin{aligned}
& V_{11}=\left(z-D_{n}\right)^{-1}\left(I d-P_{n}\right)^{-1} S^{n} \hat{\psi}_{2} \\
& V_{12}=\left(z-D_{n}\right)^{-1} R_{n}^{(2)}\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1} \\
& V_{21}=\left(z-D_{n}\right)^{-1} R_{n}^{(1)}\left(I d-P_{n}\right)^{-1} S^{n} \hat{\psi}_{2} \\
& V_{22}=\left(z-D_{n}\right)^{-1}\left(I d-Q_{n}\right)^{-1} S^{n} J \hat{\psi}_{1} .
\end{aligned}
$$

Proposition 3.1 Let $0<\delta \leq 1, M \geq 1$ and $w$ be a $\delta$-weight. Then there exist $C \equiv C(\delta, M) \geq 1$ and $N \equiv N(M, \delta) \geq 1$ such that for 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$ and $|n| \geq N$
(i) $\frac{1}{2} \leq\left|x^{F}\right|^{2}+\left|y^{F}\right|^{2} \leq 1$
(ii) $\left\|\breve{F}_{2}\right\| \leq 2 \frac{C}{\langle n\rangle^{\delta}} ; \quad\left\|J \breve{F}_{1}\right\| \leq 2 \frac{C}{\langle n\rangle^{\delta}}$
where $\|\cdot\|$ stands for the $\ell^{2}$-norm.
Proof As $\|F\|=1$, we have

$$
\|F\|^{2}=\left|x^{F}\right|^{2}+\left|y^{F}\right|^{2}+\left\|\breve{F}_{2}\right\|^{2}+\left\|J \breve{F}_{1}\right\|^{2}=1 .
$$

Hence

$$
\left|x^{F}\right|^{2}+\left|y^{F}\right|^{2} \leq 1
$$

Further, by Proposition 2.4, for $|n| \geq N_{0}$

$$
\left\|\left(I d-P_{n}\right)^{-1}\right\|_{\mathcal{L}\left(\ell_{S n_{w}}^{2}\right)} \leq 2 ;\left\|\left(I d-Q_{n}\right)^{-1}\right\|_{\mathcal{L}\left(\ell_{S}^{n_{w}}\right)}^{2} \leq 2 .
$$

By Corollary 2.3, there exists $C>1$ such that

$$
\left\|R_{n}^{(j)}\right\|_{\mathcal{L}\left(\ell_{S^{n}}^{2}, \ell^{2}\right)} \leq \frac{C}{\langle n\rangle^{\delta \wedge 1}}
$$

and by the definition of $D_{n}$, for some $1 \leq C<\infty$,

$$
\begin{aligned}
& \left\|\left(z-D_{n}\right)^{-1}\right\|_{\mathcal{L}\left(\ell_{S}^{2} n_{w}, \ell^{2}\right)} \leq \frac{C}{\langle n\rangle^{1 \wedge \delta}} \\
& \left\|\left(z-D_{n}\right)^{-1}\right\|_{\mathcal{L}\left(\ell^{2}, \ell^{2}\right)} \leq 1
\end{aligned}
$$

Hence for $|n| \geq N_{0}$

$$
\begin{aligned}
\left\|V_{11}\right\|+\left\|V_{22}\right\| & \leq \frac{C}{\langle n\rangle^{\delta \wedge 1}} \\
\left\|V_{12}\right\|+\left\|V_{21}\right\| & \leq \frac{C}{\langle n\rangle^{\delta \wedge 1}}
\end{aligned}
$$

for some $1<C<\infty$ and one concludes that

$$
\begin{aligned}
& \left\|\breve{F}_{2}\right\| \leq \frac{C}{\langle n\rangle^{\delta \wedge 1}}\left(\left|x^{F}\right|+\left|y^{F}\right|\right) \leq 2 \frac{C}{\langle n\rangle^{\delta \wedge 1}} \\
& \left\|J \breve{F}_{1}\right\| \leq \frac{C}{\langle n\rangle^{\delta \wedge 1}}\left(\left|x^{F}\right|+\left|y^{F}\right|\right) \leq 2 \frac{C}{\langle n\rangle^{\delta \wedge 1}} .
\end{aligned}
$$

By choosing $N \geq N_{0}$ sufficiently large we have, for $|n| \geq N$,

$$
\left\|\breve{F}_{2}\right\|^{2}+\left\|J \breve{F}_{1}\right\|^{2} \leq \frac{1}{2}
$$

and hence $\frac{1}{2} \leq\left|x^{F}\right|^{2}+\left|y^{F}\right|^{2}$.

### 3.2 Normal form of $L$

In this section we want to derive a normal form of the restriction of $L$ to the Riesz spaces $E_{n}$. For this purpose introduce an orthonormal basis of $E_{n}$ as follows: Choose $F \equiv F^{+}$to be an $L_{2}$-normalized eigenfunction of $L$ corresponding to the eigenvalue $\lambda^{+} \equiv \lambda_{n}^{+}$and $\Phi \in E_{n}$ with

$$
(\Phi, F)=0 ;\|\Phi\|=1
$$

where, as usual, $(\Phi, F)=\int_{0}^{1} \overline{\Phi(x)} F(x) d x$. In case $\lambda^{+}$is a double eigenvalue,

$$
\binom{L \Phi}{L F}=\left(\begin{array}{cc}
\lambda^{+} & \xi  \tag{3.1}\\
0 & \lambda^{+}
\end{array}\right)\binom{\Phi}{F}
$$

where $\xi \equiv \xi_{n}$ vanishes iff $\lambda^{+}$is of geometric multiplicity two.
In case $\lambda_{n}^{-} \neq \lambda_{n}^{+}$, choose an $L_{2}$-normalized eigenfunction $F^{-}$of $\lambda^{-} \equiv \lambda_{n}^{-}$. Then

$$
F^{-}=a F+b \Phi ;|a|^{2}+|b|^{2}=1 ; b \neq 0
$$

With $\Phi=\frac{1}{b} F^{-}-\frac{a}{b} F$,

$$
\begin{aligned}
L \Phi & =\lambda^{-} \frac{1}{b} F^{-}+\lambda^{+} \frac{a}{b} F \\
& =\lambda^{-}\left(\frac{1}{b} F^{-}-\frac{a}{b} F\right)-\gamma \frac{a}{b} F
\end{aligned}
$$

where $\gamma \equiv \gamma_{n}:=\lambda^{+}-\lambda^{-}$. Hence

$$
\binom{L \Phi}{L F}=\left(\begin{array}{cc}
\lambda^{-} & \xi  \tag{3.2}\\
0 & \lambda^{+}
\end{array}\right)\binom{\Phi}{F}
$$

with $\xi \equiv \xi_{n}:=-\gamma \frac{a}{b}$. Notice that (3.1) and (3.2) have the same form. We refer to this form as the normal form of the restriction of $L$ to the Riesz space $E_{n}$.

In the remaining part of this section we want to estimate the size of $\left(\xi_{n}\right)_{|n| \geq N}$. To this end, we write the equation $\left(L-\lambda^{-}\right) \Phi=\xi F$ in the basis $e_{k}^{+}, e_{k}^{-}(k \in \mathbb{Z})$. With $\Phi=x^{\Phi} e_{n}^{+}+y^{\Phi} e_{n}^{-}+\sum_{k \neq n} \breve{\Phi}_{2}(k) e_{k}^{+}+\breve{\Phi}_{1}(-k) e_{k}^{-}$and $F=x^{F} e_{n}^{+}+y^{F} e_{n}^{-}+$ $\sum_{k \neq n} \breve{F}_{2}(k) e_{k}^{+}+\breve{F}_{1}(-k) e_{k}^{-}$, we then obtain the following inhomogeneous system (cf. (2.8) - (2.10))

$$
\begin{align*}
& -z^{-} x^{\Phi}+\hat{\psi}_{2}(2 n) y^{\Phi}+\left\langle S^{n} \hat{\psi}_{2}, J \breve{\Phi}_{1}\right\rangle=\xi x^{F}  \tag{3.3}\\
& \hat{\psi}_{1}(-2 n) x^{\Phi}-z^{-} y^{\Phi}+\left\langle S^{n} J \hat{\psi}_{1}, \breve{\Phi}_{2}\right\rangle=\xi y^{F}  \tag{3.4}\\
& \binom{y^{\Phi}\left(S^{n} \hat{\psi}_{2}\right)_{\mathbb{Z} \backslash n}}{x^{\Phi}\left(S^{n} J \hat{\psi}_{1}\right)_{\mathbb{Z} \backslash n}}+\left(A_{n}-z^{-}\right)\binom{\breve{\Phi}_{2}}{J \breve{\Phi}_{1}}=\xi\binom{\breve{F}_{2}}{J \breve{F}_{1}} \tag{3.5}
\end{align*}
$$

where, as usual, $\lambda_{n}^{-} \equiv \lambda^{-}=n \pi+z^{-}$. We use the above system to obtain an estimate for $\xi \equiv \xi_{n}$.
Write $\breve{\Phi}=\left(\breve{\Phi}_{2}, J \breve{\Phi}_{1}\right)$ and $\breve{F}=\left(\breve{F}_{2}, J \breve{F}_{1}\right)$. Recall that $w$ is assumed to be a $\delta$-weight and hence by Corollary 2.5, equation (3.5) (with $|n| \geq N_{0}$ ) can be solved for $\breve{\Phi}$,

$$
\breve{\Phi}=\left(z^{-}-A_{n}\right)^{-1}\binom{y^{\Phi}\left(S^{n} \hat{\psi}_{2}\right)_{\mathbb{Z} \backslash n}}{x^{\Phi}\left(S^{n} J \hat{\psi}_{1}\right)_{\mathbb{Z} \backslash n}}-\xi\left(z^{-}-A_{n}\right)^{-1} \breve{F}
$$

In this form, $\breve{\Phi}$ is substituted into (3.3) - (3.4) to obtain (cf. Corollary 2.5)

$$
\begin{align*}
& \left(\begin{array}{cc}
-z^{-}+\alpha\left(n, z^{-}\right) & \hat{\psi}_{2}(2 n)+\beta^{+}\left(n, z^{-}\right) \\
\hat{\psi}_{1}(-2 n)+\beta^{-}\left(n, z^{-}\right) & -z^{-}+\alpha\left(n, z^{-}\right)
\end{array}\right)\binom{x^{\Phi}}{y^{\Phi}}  \tag{3.6}\\
& =\xi\binom{x^{F}}{y^{F}}+\xi\left\langle\binom{ S^{n} \hat{\psi}_{2}}{S^{n} J \hat{\psi}_{1}},\left(z^{-}-A_{n}\right)^{-1} \breve{F}\right\rangle .
\end{align*}
$$

Denote the right side of (3.6) by $R S$. By Corollary $2.5\left(z^{-}-A_{n}\right)^{-1}$ is uniformly bounded for $|n|$ sufficiently large and by Proposition 3.1, for $|n| \geq$ $N$,

$$
\frac{1}{2} \leq\left|x^{F}\right|^{2}+\left|y^{F}\right|^{2} ;\|\breve{F}\| \leq \frac{C}{\langle n\rangle^{\delta}}
$$

Hence $R S$ can be estimated from below: There exists $1 \leq C \equiv C_{\delta, M}<\infty$ so that for $|n| \geq N(N$ as in Proposition 3.1)

$$
R S \geq|\xi|\left(\frac{1}{\sqrt{2}}-\frac{C}{\langle n\rangle^{\delta}}\right)
$$

By choosing $N$ larger if necessary, we can assume that

$$
\begin{equation*}
\frac{1}{\sqrt{2}}-\frac{C}{\langle n\rangle^{\delta}} \geq \frac{1}{2} \quad \forall|n| \geq N \tag{3.7}
\end{equation*}
$$

and (3.6) leads to

$$
\begin{equation*}
\left|\xi_{n}\right| \leq 4\left(\left|\zeta_{n}^{-}\right|+\left|\hat{\psi}_{1}(-2 n)\right|+\left|\hat{\psi}_{2}(2 n)\right|+\left|\beta^{+}\left(n, z^{-}\right)\right|+\left|\beta^{-}\left(n, z^{-}\right)\right|\right) \tag{3.8}
\end{equation*}
$$

where we used that $\left|x^{\Phi}\right|^{2}+\left|y^{\Phi}\right|^{2} \leq 1$ and $\zeta_{n}^{-}=z^{-}-\alpha\left(n, z^{-}\right)$with $z^{-} \equiv z\left(\zeta_{n}^{-}\right)$.
In view of Proposition 2.10 and Lemma 2.13, one then concludes from (3.8) the following

Proposition 3.2 Let $M \geq 1,0<\delta$, and $w$ be a $\delta$-weight. Then there exist $1 \leq N<\infty, 1 \leq C=C_{\delta}<\infty$ such that for any 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$

$$
\sum_{|n| \geq N} w(2 n)^{2}\left|\xi_{n}\right|^{2} \leq C
$$

## 4 Dirichlet eigenvalues

### 4.1 Dirichlet boundary value problem

Consider the Zakharov-Shabat operator $L \equiv L\left(\psi_{1}, \psi_{2}\right)$ on the interval $[0,1]$.
Definition 4.1 $F=\left(F_{1}, F_{2}\right) \in H^{1}\left([0,1] ; \mathbb{C}^{2}\right)$ satisfies Dirichlet boundary conditions if

$$
\begin{equation*}
F_{1}(0)-F_{2}(0)=0 ; F_{1}(1)-F_{2}(1)=0 \tag{4.1}
\end{equation*}
$$

We mention that the Dirichlet boundary conditions take a more familiar form when the operator $L$ is written as an $A K N S$ operator $L_{A K N S}$

$$
L_{A K N S}=\left(\begin{array}{cc}
0 & -1  \tag{4.2}\\
1 & 0
\end{array}\right) \frac{d}{d x}+\left(\begin{array}{cc}
-q & p \\
p & q
\end{array}\right)
$$

where $\left(\psi_{1}, \psi_{2}\right)$ and $(p, q)$ are related by

$$
\psi_{1}=-q+i p ; \psi_{2}=-q-i p
$$

If $F=\left(F_{1}, F_{2}\right) \in H^{1}\left([0,1] ; \mathbb{C}^{2}\right)$ satisfies $L F=\lambda F$, then $L \tilde{F}=\lambda \tilde{F}$ where $\tilde{F}=\left(\tilde{F}_{1}, \tilde{F}_{2}\right)$ is given by

$$
\tilde{F}_{1}=\frac{1}{\sqrt{2} i}\left(F_{1}+F_{2}\right) ; \quad \tilde{F}_{2}=\frac{1}{\sqrt{2}}\left(F_{2}-F_{1}\right) .
$$

The Dirichlet boundary conditions (4.1) then take the familiar form

$$
\tilde{F}_{2}(0)=0 ; \quad \tilde{F}_{2}(1)=0
$$

For the remaining part of section 4 , let $M \geq 1, \delta>0$, and a $\delta$-weight $w$ be given as well as arbitrary 1-periodic functions $\psi_{1}, \psi_{2} \in H^{w}$ with $\left\|\psi_{j}\right\|_{w}<M$. In Appendix A we have introduced, for $|n| \geq N$ with $N$ given by Theorem 1.1, the Riesz projectors $\Pi_{2 n}, \Pi_{2 n-1}$ corresponding to periodic resp. antiperiodic boundary value problem on $[0,1]$ for $L$ and the two dimensional subspaces $E_{n}=\operatorname{Range}\left(\Pi_{n}\right)$.
The following proposition assures that there exists a 1-dimensional subspace of $E_{n}$ which satisfies Dirichlet boundary conditions. Let $(F, \Phi)$ denote the orthonormal basis of $E_{n} \subseteq L^{2}\left([0,1] ; \mathbb{C}^{2}\right)$, introduced in section 3.2.

Proposition 4.2 For any $|n| \geq N$, there exists $G=\left(G_{1}, G_{2}\right) \in E_{n}$

$$
G=\alpha F+\beta \Phi ;|\alpha|^{2}+|\beta|^{2}=1
$$

which satisfies Dirichlet boundary conditions

$$
G_{1}(0)-G_{2}(0)=0 ; \quad G_{1}(1)-G_{2}(1)=0
$$

Proof First consider the case where $F$ satisfies $F_{1}(0)-F_{2}(0)=0$. As $F$ is either periodic or antiperiodic we conclude that $F_{1}(1)-F_{2}(1)=0$ as well and thus $G:=F$ has the required properties. If $F_{1}(0)-F_{2}(0) \neq 0$, notice that

$$
\tilde{G}(x):=\left(F_{1}(0)-F_{2}(0)\right) \Phi(x)-\left(\Phi_{1}(0)-\Phi_{2}(0)\right) F(x)
$$

satisfies Dirichlet boundary conditions. As $\tilde{G} \not \equiv 0$, we may define

$$
G:=\frac{\tilde{G}}{\|\tilde{G}\|}
$$

By (3.1) - (3.2), $L \Phi=\lambda^{-} \Phi+\xi F$ and $L F=\lambda^{+} F$, hence, with $\gamma \equiv \gamma_{n}=$ $\lambda^{+}-\lambda^{-}$and $\lambda \equiv \lambda^{+}$

$$
\begin{align*}
L G & =\alpha \lambda F+\beta L \Phi \\
& =\lambda G-\beta \gamma \Phi+\beta \xi F . \tag{4.3}
\end{align*}
$$

For $|n| \geq N$ sufficiently large, $\xi \equiv \xi_{n}$ and $\gamma \equiv \gamma_{n}$ are small and $G$ is almost a Dirichlet eigenfunction. In the next sections we prove that $\lambda \equiv \lambda_{n}^{+}$and $G$ are good approximations of the Dirichlet eigenvalue $\mu \equiv \mu_{n}$ respectively Dirichlet eigenfunction $H$.

### 4.2 Decomposition

Let $L_{D i r}$ denote the closed operator $L_{D i r}=L\left(\psi_{1}, \psi_{2}\right)$ with domain

$$
\operatorname{dom} L_{\text {Dir }}:=\left\{F \in H^{1}[0,1] \mid F_{1}(0)-F_{2}(0)=0 ; F_{1}(1)-F_{2}(1)=0\right\} .
$$

Let us fix $n$ with $|n| \geq N$ ( $N$ as in Theorem 1.1). $\Pi_{D i r} \equiv \Pi_{n, D i r}$ denotes the Riesz projector

$$
\Pi_{D i r}:=\frac{1}{2 \pi i} \int_{|z-n \pi|=\frac{\pi}{2}}\left(z-L_{D i r}\right)^{-1} d z
$$

acting on $L^{2}\left([0,1] ; \mathbb{C}^{2}\right)$ (cf. Appendix A). Let $\Omega_{D i r}:=I d-\Pi_{D i r}$.
Notice that

$$
\text { Range } \Pi_{\text {Dir }}=\{a H \mid a \in \mathbb{C}\}
$$

where $H \in \operatorname{dom} L_{D i r}$ is an $L^{2}[0,1]$-normalized eigenfunction for the Dirichlet eigenvalue $\mu \equiv \mu_{n}$,

$$
L_{D i r} H=\mu H ;\|H\|=1
$$

Let $\chi \in \mathbb{C}$ with $|\chi| \leq 1$ defined by $\Pi_{D i r} G=\chi H$ where $G$ is given by Proposition 4.2. We have

$$
G=\chi H+\Omega_{D i r} G
$$

As $G$ and $H$ are in $d o m L_{D i r}, \Omega_{D i r} G \in d o m L_{D i r}$ and

$$
\begin{equation*}
L_{D i r} G=\chi \mu H+L_{D i r} \Omega_{D i r} G=\chi \mu H+\Omega_{D i r} L_{D i r} \Omega_{D i r} G . \tag{4.4}
\end{equation*}
$$

Where for the last equality we have used, $\Pi_{D i r} L_{D i r} \Omega_{D i r} G=0$, as $L_{D i r}$ and $\Pi_{D i r}$ commute on $d o m L_{D i r}$ and $\Pi_{D i r} \Omega_{D i r}=0$.
On the other hand by (4.3),

$$
L G=\lambda G+R ; \quad R=-\beta \gamma \Phi+\beta \xi F
$$

and thus, with $G=\chi H+\Omega_{D i r} G$,

$$
\begin{equation*}
L_{D i r} G=\lambda \chi H+\lambda \Omega_{D i r} G+\left(\Pi_{D i r}+\Omega_{D i r}\right) R \tag{4.5}
\end{equation*}
$$

Comparing the decompositions of the right sides of (4.4) and (4.5) leads to the following

## Lemma 4.3

$$
\begin{align*}
\chi(\mu-\lambda) H & =\Pi_{D i r} R ;  \tag{4.6}\\
\left(L_{D i r}-\lambda\right)\left(\Omega_{D i r} G\right) & =\Omega_{D i r} R \tag{4.7}
\end{align*}
$$

where $R$ is given by

$$
\begin{equation*}
R=-\beta \gamma \Phi+\beta \xi F \tag{4.8}
\end{equation*}
$$

### 4.3 Proof of Theorem 1.2

The equations (4.6) - (4.8) are now used to obtain estimates for $\mid \mu_{n}-$ $\lambda_{n}^{+} \mid(|n| \geq N)$. For this we need to establish that $|\chi| \leq 1$ is bounded away from 0 and that $\left\|\Pi_{D i r} R\right\|$ is small. The latter is easily seen as $\|R\| \leq|\gamma|+|\xi|$. To verify that $|\chi|$ is bounded away from 0 we show that $\Omega_{D i r} G=G-\chi G$ is small. This is proved by using equation (4.7).

Lemma 4.4 There exists $N \geq 1$ so that

$$
\left|\chi_{n}\right| \geq \frac{1}{2} \quad \forall|n| \geq N
$$

Proof As $G=\chi H+\Omega_{D i r} G$,

$$
|\chi|\|H\|=\|G\|-\left\|\Omega_{D i r} G\right\|=1-\left\|\Omega_{D i r} G\right\| .
$$

By Lemma 4.3 and Lemma A. 2 (for (4.9)), Proposition 3.2 (for (4.10)), and Theorem 1.1 (i) (for (4.11)) there exist $1 \leq N<\infty$ and $1 \leq C<\infty$ so that for $|n| \geq N$

$$
\begin{gather*}
\left\|\Omega_{D i r} G\right\|=\left\|\left(L_{D i r}-\lambda\right)^{-1}\left(\Omega_{D i r} R\right)\right\| \leq C\|R\| \leq C\left(\left|\xi_{n}\right|+\left|\gamma_{n}\right|\right)  \tag{4.9}\\
\left|\xi_{n}\right| \leq \frac{C}{\langle n\rangle^{\delta}}  \tag{4.10}\\
\left|\gamma_{n}\right| \leq \frac{C}{\langle n\rangle^{\delta}}, \tag{4.11}
\end{gather*}
$$

(where for the last two inequalities we used that $w$ is a $\delta$-weight). Combining the above inequalities shows that for $|n|$ large enough

$$
\left|\chi_{n}\right| \geq \frac{1}{2}
$$

Proof of Theorem 1.2 By (4.6),

$$
\left|\chi\|\mu-\lambda \mid\| H\|=\| \Pi_{D i r} R \| .\right.
$$

By Lemma 4.4 there exists $N \geq 1$ so that for $|n| \geq N$

$$
\left|\mu_{n}-\lambda_{n}^{+}\right| \leq 2 C\left(\left|\xi_{n}\right|+\left|\gamma_{n}\right|\right)
$$

where we have used that $\left\|\Pi_{D i r}\right\| \leq C$ (cf. Lemma A.2). The claimed estimate then follows from the estimates of $\xi_{n}$ (Proposition 3.2) and of $\gamma_{n}$ (Theorem 1.1 (i)).

## A Appendix A: Spectral properties of $L\left(\psi_{1}, \psi_{2}\right)$

In this appendix we consider the operator $L\left(\psi_{1}, \psi_{2}\right)\left(\psi_{1}, \psi_{2} 1\right.$-periodic functions in $\left.L^{2}\left([0,2], \mathbb{C}^{2}\right)\right)$ with various boundary conditions. For $b c \in\left\{\right.$ Dir, Per $\left.{ }^{ \pm}, \operatorname{Per}\right\}$
denote by $L_{b c}$ the Zakharov-Shabat operator $L_{b c}=L\left(\psi_{1}, \psi_{2}\right)$ with the following domains:

$$
\begin{aligned}
\operatorname{dom} L_{\text {Dir }} & :=\left\{F \in H^{1}[0,1] \mid F_{1}(0)-F_{2}(0)=0 ; F_{1}(1)-F_{2}(1)=0\right\} ; \\
\operatorname{dom} L_{\text {Per}}+ & :=\left\{F \in H^{1}[0,1] \mid F(0)=F(1)\right\} ; \\
\operatorname{dom} L_{P e r^{-}} & :=\left\{F \in H^{1}[0,1] \mid F(0)=-F(1)\right\} .
\end{aligned}
$$

The operator $L \equiv L_{P e r}$ is defined on the interval $[0,2]$ and has the following domain,

$$
\operatorname{dom} L_{P e r}:=\left\{F \in H^{1}[0,2] \mid F(0)=F(2)\right\} .
$$

Let $\operatorname{spec}_{b c} \equiv \operatorname{spec}\left(L_{b c}\right)$ be the spectrum of $L_{b c}$. For potentials $\psi_{1}=\psi_{2} \equiv 0$, i.e. $L_{0}:=L(0,0)=i\left(\begin{array}{cc}1 & 0 \\ 0 & -1\end{array}\right) \frac{d}{d x}, \operatorname{spec}_{b c}\left(L_{0}\right)$ can be given explicitely:

$$
\begin{align*}
\operatorname{spec}_{\text {Dir }}\left(L_{0}\right) & =\{k \pi \mid k \in \mathbb{Z}\} ;  \tag{A.1}\\
\operatorname{spec}_{\text {Per }}\left(L_{0}\right) & =\{2 k \pi \mid k \in \mathbb{Z}\} ;  \tag{A.2}\\
\operatorname{spec}_{\text {Per }}\left(L_{0}\right) & =\{2(k+1) \pi \mid k \in \mathbb{Z}\} ;  \tag{A.3}\\
\operatorname{spec}_{\text {Per }}\left(L_{0}\right) & =\{k \pi \mid k \in \mathbb{Z}\} . \tag{A.4}
\end{align*}
$$

Proposition A. 1 Let $\delta>0, M \geq 1$ and $w$ be a $\delta$-weight. There exists an even integer $N$ such that for any 1-periodic functions $\psi_{1}$ and $\psi_{2}$ in $H^{w},\left\|\psi_{j}\right\|_{w} \leq M$, the following statements hold:
(i) for $b c \in\left\{D i r, P e r r^{ \pm}, P e r\right\}$,
$\operatorname{spec}_{b c} \subset\{\lambda \in \mathbb{C}| | \lambda \mid<N \pi-\pi / 2\} \cup\left(\bigcup_{|k| \geq N}\{\lambda \in \mathbb{C}| | \lambda-k \pi \mid<\pi / 2\}\right) ;$
(ii) for $|k| \geq N$, $\operatorname{spec}_{P e r} \cap\{\lambda \in \mathbb{C}| | \lambda-k \pi \mid<\pi / 2\}$ contains exactly one isolated pair of eigenvalues;
(iii) for $|k| \geq N$ and bc $:=\operatorname{Per}^{+}(k$ even $)$ and bc $:=\operatorname{Per}^{-}$( $k$ odd), spec $_{b c} \cap\{\lambda \in \mathbb{C}| | \lambda-k \pi \mid<\pi / 2\}$ contains exactly one isolated pair of eigenvalues;
(iv) for $|k| \geq N$, spec $_{\text {Dir }} \cap\{\lambda \in \mathbb{C}| | \lambda-k \pi \mid<\pi / 2\}$ contains exactly one eigenvalue;
(v) the cardinality $N_{b c}$ of spec $_{b c} \cap\{\lambda \in \mathbb{C}| | \lambda \mid<N \pi-\pi / 2\}$ is equal to $4 N-2$ for $b c=\operatorname{Per}, 2 N-1$ for $b c=\operatorname{Dir}, 2 N-2$ for $b c=P e r^{+}$and $2 N$ for $b c=$ Per $^{-}$.

As $\operatorname{spec}_{P e r^{+}} \cup \operatorname{spec}_{P e r^{-}} \subseteq \operatorname{spec}_{P e r}$, Proposition A. 1 implies

$$
\operatorname{spec}_{P e r}=\operatorname{spec}_{P e r^{+}} \cup \operatorname{spec}_{P e r^{-}} .
$$

Proof Define for $n \geq 1$, the union of contours,

$$
\mathcal{R}_{n}=\{\lambda \in \mathbb{C}| | \lambda \mid=n \pi-\pi / 2\} \cup\left(\bigcup_{|k|>n}\{\lambda \in \mathbb{C}| | \lambda-k \pi \mid=\pi / 2\}\right)
$$

By (A.1) - (A.4), $\left(L_{0}-\lambda\right): \operatorname{dom}\left(L_{b c}\right) \rightarrow L^{2}$ is invertible for any $\lambda \in \mathcal{R}_{n}$, hence

$$
\begin{equation*}
(L-\lambda)=\left(L_{0}-\lambda\right)\left(I d+Q_{\lambda}\right) \tag{A.5}
\end{equation*}
$$

where

$$
Q_{\lambda}=\left(L_{0}-\lambda\right)^{-1}\left(\begin{array}{cc}
0 & \psi_{1} \\
\psi_{2} & 0
\end{array}\right)
$$

Using the orthogonal decomposition of $L^{2}$ by the eigenfunctions of $\left(L_{0}\right)_{b c}$ and the assumption that $w$ is a $\delta$-weight, one gets (with $\mathcal{L} \equiv \mathcal{L}\left(L^{2}\right)$ )

$$
\begin{equation*}
\left\|Q_{\lambda}\right\|_{\mathcal{L}} \leq M\left(\sum_{k \in \mathbb{Z}} \frac{1}{|k|^{2 \delta}|k \pi-\lambda|^{2}}\right)^{1 / 2} \tag{A.6}
\end{equation*}
$$

As $\max _{k \in \mathbb{Z}}\left(\frac{1}{|k|^{2 \delta}\langle k-n\rangle^{1 / 2}}\right)^{1 / 2} \leq \frac{1}{\langle n\rangle^{\delta \Lambda 1 / 4}}$, one deduces from (A.6) that, for $\lambda \in$ $\mathcal{R}_{n}$,

$$
\begin{equation*}
\left\|Q_{\lambda}\right\|_{\mathcal{L}} \leq \frac{M}{\langle n\rangle^{\delta \wedge 1 / 4}}\left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k\rangle^{3 / 2}}\right)^{1 / 2} \tag{A.7}
\end{equation*}
$$

Let $N$ be an even integer such that

$$
\frac{M}{\langle n\rangle^{\delta \wedge 1 / 4}}\left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k\rangle^{3 / 2}}\right)^{1 / 2} \leq 1 / 2
$$

Then, for $\lambda \in \mathcal{R}_{n}$ with $n \geq N$

$$
\begin{equation*}
\left\|Q_{\lambda}\right\|_{\mathcal{L}} \leq 1 / 2 \tag{A.8}
\end{equation*}
$$

Combining (A.5) and (A.8), one deduces that $(L-\lambda): \operatorname{dom}\left(L_{b c}\right) \rightarrow L^{2}$ is invertible for any $\lambda \in \mathcal{R}_{n}(n \geq N)$ and any 1-periodic functions $\psi_{1}, \psi_{2}$ in $H^{w}$ with $\left\|\psi_{j}\right\|_{w} \leq M$. In particular, $\mathcal{R}_{n}(n \geq N)$ is contained in the resolvent set of $L_{b c}\left(t \psi_{1}, t \psi_{2}\right)$ for any $0 \leq t \leq 1$. Hence the number of eigenvalues of $L_{b c}\left(t \psi_{1}, t \psi_{2}\right)$ in each connected component of the interior of $\mathcal{R}_{n}$ stays the same for any $0 \leq t \leq 1$. To see that all eigenvalues are inside $\mathcal{R}_{n}$ one chooses $n$ bigger and bigger.

It follows from Proposition A. 1 that the Riesz projectors $\Pi_{n}$ and $\Pi_{n, \text { Dir }}$ are well defined (for any $|n| \geq N$ and 1-periodic functions $\psi_{1}, \psi_{2}$ with $\left\|\psi_{j}\right\|_{w} \leq$ M)

$$
\begin{aligned}
& \Pi_{n}:=\frac{1}{2 \pi i} \int_{|\lambda-n \pi|=\pi / 2}\left(z-L_{P e r^{+}}\right)^{-1} d z \quad(n \text { even },|n| \geq N) \\
& \Pi_{n}:=\frac{1}{2 \pi i} \int_{|\lambda-n \pi|=\pi / 2}\left(z-L_{P e r^{-}}\right)^{-1} d z \quad(n \text { odd },|n| \geq N)
\end{aligned}
$$

and

$$
\Pi_{n, \text { Dir }}:=\frac{1}{2 \pi i} \int_{|\lambda-n \pi|=\pi / 2}\left(z-L_{D i r}\right)^{-1} d z(|n| \geq N),
$$

where the contours $\{\lambda||\lambda-n \pi|=\pi / 2\}$ in the integrals above are counterclockwise oriented. Furthermore, using (A.5) and (A.8), one deduces

Lemma A. 2 Assume that the assumptions of Proposition A. 1 hold. Then there exists a constant $1 \leq C \leq \infty$ such that for any $|n| \geq N$ (with $N$ as in Proposition A.1)

$$
\left\|\Pi_{n}\right\|_{\mathcal{L}\left(L^{2}[0,1]\right)} \leq C
$$

and

$$
\left\|\Pi_{n, D i r}\right\|_{\mathcal{L}\left(L^{2}[0,1]\right)} \leq C
$$

## B Appendix B: Proof of Lemma 2.8

W.l.o.g. we may assume that $\delta_{*}=\delta$.
(i) As $w_{*}$ is submultiplicative, one has

$$
\begin{equation*}
w(2 n)=\langle 2 n\rangle^{\delta} w_{*}(2 n) \leq 2^{\delta}\langle n\rangle^{\delta} w_{*}(n+k) w_{*}(k+j) w_{*}(j+n) \tag{B.1}
\end{equation*}
$$

and, by assumption, $\langle n\rangle^{\alpha} d_{n}(k) \leq d(k)(\forall k)$. This leads to

$$
\begin{align*}
& \langle n\rangle^{2 \delta+\alpha} w(2 n) \Psi_{n}\left(a, b, d_{n}\right)(2 n) \\
& \quad \leq\langle n\rangle^{2 \delta} w(2 n) \sum_{k} \frac{a(k+n)}{\langle k-n\rangle} \sum_{j} \frac{b(k+j)}{\langle j-n\rangle} d(j+n)  \tag{B.2}\\
& \quad \leq \sum_{k, j} K_{n}(k, j) \tilde{a}(k+n) \tilde{b}(k, j) \tilde{d}(j+n)
\end{align*}
$$

where for any $u \in \ell_{w}^{2}$ we denote by $\tilde{u}$ the $\ell^{2}$-sequence $\tilde{u}(j):=w(j) u(j)$ and $K_{n}(k, j)$ is given by

$$
K_{n}(k, j):=\frac{2^{\delta}\langle n\rangle^{3 \delta}}{\langle k-n\rangle\langle j-n\rangle\langle k+n\rangle^{\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} .
$$

Notice that $K_{n}(k, j)$ is symmetric in $k$ and $j$. To estimate $K_{n}(k, j)$ we need to consider four different regions:
Estimate of $K_{n}(k, j)$ in $|k-n|<\frac{|n|}{2},|j-n|<\frac{|n|}{2}$ : In this case

$$
|k+n| \geq|n| ;|j+n| \geq|n| ;|k+j| \geq 2|n|-|k-n|-|j-n| \geq|n|
$$

hence

$$
K_{n}(k, j) \leq \frac{2^{\delta}}{\langle k-n\rangle\langle j-n\rangle} \leq \frac{1}{\langle k-n\rangle^{2}}+\frac{1}{\langle j-n\rangle^{2}}
$$

Estimate of $K_{n}(k, j)$ in $|k-n| \geq \frac{|n|}{2},|j-n|<\frac{|n|}{2}$ : In this case

$$
|k-n| \geq \frac{|n|}{2} ;|j+n|>|n|
$$

hence

$$
K_{n}(k, j) \leq \frac{2^{\delta}}{\langle j-n\rangle\langle k+j\rangle^{\delta}}
$$

Estimate of $K_{n}(k, j)$ in $|k-n|<\frac{|n|}{2},|j-n| \geq \frac{|n|}{2}$ : Using the symmetry of $K_{n}(k, j)$ in $k$ and $j$, the latter estimate leads to

$$
K_{n}(k, j) \leq \frac{2^{\delta}}{\langle k-n\rangle\langle k+j\rangle^{\delta}}
$$

Estimate of $K_{n}(k, j)$ in $|k-n| \geq \frac{|n|}{2},|j-n| \geq \frac{|n|}{2}$ : We get

$$
K_{n}(k, j) \leq \frac{16^{\delta}}{\langle k-n\rangle^{1-\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} .
$$

Combining the above estimates one obtains for $k, j, n \in \mathbb{Z}$,
$K_{n}(k, j) \leq \frac{1}{\langle k-n\rangle^{2}}+\frac{1}{\langle j-n\rangle^{2}}+\frac{2}{\langle k+j\rangle^{\delta}} \frac{1}{\langle k-n\rangle}+\frac{4}{\langle k-n\rangle^{1-\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}}$.
Therefore

$$
\begin{align*}
& \sum_{k, j} K_{n}(k, j) \tilde{a}(k+n) \tilde{b}(k+j) \tilde{d}(j+n) \\
& \leq\left(\tilde{a} * \frac{1}{\langle k\rangle^{2}}(J \tilde{b} * \tilde{d})\right)(2 n)+\left(\tilde{d} * \frac{1}{\langle k\rangle^{2}}(J \tilde{b} * \tilde{a})\right)(2 n)+  \tag{B.3}\\
& +2\left(\tilde{a} * \frac{1}{\langle k\rangle}\left(\frac{J \tilde{b}}{\langle k\rangle^{\delta}} * \tilde{d}\right)\right)(2 n)+4\left(\tilde{a} * \frac{1}{\langle k\rangle^{1-\delta}}\left(\frac{J \tilde{b}}{\langle k\rangle^{\delta}} * \frac{\tilde{d}}{\langle k\rangle^{\delta}}\right)\right)
\end{align*}
$$

where for $u \in \ell^{2}(\mathbb{Z})$ and $\eta \geq 0, \frac{u}{\langle k\rangle^{\eta}}$ denotes the sequence given by $\left(\frac{u}{\langle k\rangle^{\eta}}\right)(j):=$ $\frac{u(j)}{\langle j\rangle^{\eta}}(\forall j)$. Using the standard convolution estimates $\|u * v\|_{\ell^{2}} \leq\|u\|_{\ell^{1}}\|v\|_{\ell^{2}}$ and $\|u * v\|_{\ell^{\infty}} \leq\|u\|_{\ell^{2}}\|v\|_{\ell^{2}}$ for the first two terms on the right side of (B.3), Corollary B. 2 (i) for the third term and Corollary B. 2 (ii) for the last term on the right side of (B.3), one obtains from (B.2)

$$
\sum_{n}\left(\langle n\rangle^{2 \delta+\alpha} w(2 n) \Psi_{n}\left(a, b, d_{n}\right)(2 n)\right)^{2} \leq C\|a\|_{w}\|b\|_{w}\|d\|_{w}
$$

for a constant $1 \leq C \leq C_{\delta}<\infty$ only depending on $\delta$.
(ii) Using (B.1) and the assumption $\langle n\rangle^{\alpha} d_{n}(k) \leq d(k)(\forall k)$ we get

$$
\begin{aligned}
& \langle n\rangle^{\delta+\alpha} w(n+\ell) \Psi_{n}\left(a, b, d_{n}\right)(\ell+n) \leq \\
& \quad \leq\langle n\rangle^{\delta} w(n+\ell) \sum_{k, j} \frac{a(k+\ell)}{\langle k-n\rangle} \frac{b(k+j)}{\langle k-j\rangle} d(j+n) \\
& \quad \leq \sum_{k, j} H_{n}(\ell, k, j) \tilde{a}(k+\ell) \tilde{b}(k+j) \tilde{d}(j+n)
\end{aligned}
$$

where $H_{n}(\ell, k, j)$ is given by

$$
H_{n}(\ell, k, j):=\frac{\langle n\rangle^{\delta}\langle\ell+n\rangle^{\delta}}{\langle k-n\rangle\langle j-n\rangle\langle k+\ell\rangle^{\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} .
$$

To estimate $H_{n}(\ell, k, j)$ we need to consider two different regions:
Estimate of $H_{n}(\ell, k, j)$ in $|j-n|<\frac{|n|}{2}$ : In this case

$$
|j+n|>|n| ;\langle\ell+n\rangle^{\delta} \leq\langle\ell+k\rangle^{\delta}\langle-k+n\rangle^{\delta},
$$

hence

$$
\begin{aligned}
H_{n}(\ell, k, j) & \leq \frac{1}{\langle k-n\rangle^{1-\delta}\langle j-n\rangle\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} \\
& \leq \frac{1}{\langle k-n\rangle^{1-\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} .
\end{aligned}
$$

Estimate of $H_{n}(\ell, k, j)$ in $|j-n|>\frac{|n|}{2}$ : In this case

$$
2|j-n|>|n| ;\langle\ell+n\rangle^{\delta} \leq\langle\ell+k\rangle^{\delta}\langle-k+n\rangle^{\delta},
$$

hence

$$
\begin{aligned}
H_{n}(\ell, k, j) & \leq \frac{2^{\delta}}{\langle k-n\rangle^{1-\delta}\langle j-n\rangle^{1-\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} \\
& \leq \frac{2^{\delta}}{\langle k-n\rangle^{1-\delta}\langle k+j\rangle^{\delta}\langle j+n\rangle^{\delta}} .
\end{aligned}
$$

Hence in both cases we obtain the same estimate. Define $\tilde{e}(\ell+n) \equiv w(\ell+$ $n) e(\ell+n)$ by

$$
\tilde{e}(\ell+n):=\sum_{k, j}\left(\frac{1}{\langle k-n\rangle\langle j+n\rangle^{\delta}\langle k+j\rangle^{\delta}}\right) \tilde{a}(k+\ell) \tilde{b}(\ell+j) \tilde{d}(j+n) .
$$

Then we have

$$
\langle n\rangle^{\delta+\alpha} w(n+\ell) \Psi_{n}\left(a, b, d_{n}\right)(\ell+n) \leq w(\ell+n) e(\ell+n)
$$

and

$$
\tilde{e}(\ell)=\left(\tilde{a} * \frac{1}{\langle k\rangle^{1-\delta}}\left(\frac{J \tilde{b}}{\langle k\rangle^{\delta}} * \frac{\tilde{d}}{\langle k\rangle^{\delta}}\right)\right)(\ell) .
$$

By Corollary B. 2 (ii),

$$
\|\tilde{e}\|_{\ell^{2}} \leq C\|a\|_{w}\|b\|_{w}\|d\|_{w}
$$

for some constant $1 \leq C \equiv C_{\delta}<\infty$.
It remains to establish the auxilary results used in the proof of Lemma 2.8. First we need the following

Lemma B. 1 Let $0<\eta \leq 1$. Then
(i) $\left\|\frac{a}{\langle k\rangle^{\eta}}\right\|_{\ell^{p}} \leq C_{p, \eta}\|a\|_{\ell^{2}} \quad \forall a \in \ell^{2}$ and $p>\frac{2}{2 \eta+1}$
(ii) $\left\|\frac{a}{\langle k\rangle^{\eta}}\right\|_{\ell^{1}} \leq C_{q, \eta}\|a\|_{\ell^{q}} \quad \forall a \in \ell^{q}$ with $1 \leq q<\frac{1}{1-\eta}$.

Proof (i) follows from Hölder's inequality with $\alpha=\frac{2}{p}$ and $\beta=\frac{2}{2-p}$,

$$
\begin{aligned}
\left(\sum_{k}\left(\frac{a(k)}{\langle k\rangle^{\eta}}\right)^{p}\right)^{1 / p} & =\left(\sum_{k} a(k)^{p} \frac{1}{\langle k\rangle^{\eta p}}\right)^{1 / p} \\
& \leq\left(\sum_{\left.|a(k)|^{2}\right)^{1 / 2}\left(\sum_{k} \frac{1}{\langle k\rangle^{\eta p \beta}}\right)^{1 / \beta p}}\right.
\end{aligned}
$$

where $\eta p \beta=\eta p \frac{2}{2-p}>1$ or $2 \eta p>2-p$ as $(2 \eta+1) p>2$ by assumption.
(ii) follows from Hölder's inequality with $\alpha=q$ and $\frac{1}{\beta}=1-\frac{1}{q}=\frac{q-1}{q}$

$$
\sum_{k} \frac{|a(k)|}{\langle k\rangle^{\eta}} \leq\left(\sum_{k} a(k)^{q}\right)^{1 / q}\left(\sum_{k}\left(\frac{1}{\langle k\rangle^{\eta}}\right)^{\frac{q}{q-1}}\right)^{\frac{q-1}{q}}
$$

where $\eta \frac{q}{q-1}>1$ or $\eta q>q-1$ as $1>(1-\eta) q$ by assumption.

Recall Young's inequality

$$
\|u * v\|_{q} \leq C_{r, p, q}\|u\|_{p}\|v\|_{r}
$$

where $r, p, q \geq 1$ with $\frac{1}{p}+\frac{1}{r}=1+\frac{1}{q}$.
Corollary B. 2 Let $0<\delta \leq \frac{1}{2}$
(i) $\left\|\frac{1}{\langle k\rangle}\left(\frac{a}{\langle k\rangle^{\delta}} * b\right)\right\|_{\ell^{1}} \leq C_{\delta}\|a\|_{\ell^{2}}\|b\|_{\ell^{2}} \quad \forall a, b \in \ell^{2}$
(ii) $\left\|\frac{1}{\langle k\rangle^{1-\delta}}\left(\frac{a}{\langle k\rangle^{\delta}} * \frac{b}{\langle k\rangle^{\delta}}\right)\right\|_{\ell^{1}} \leq C_{\delta}\|a\|_{\ell^{2}}\|b\|_{\ell^{2}} \quad \forall a, b \in \ell^{2}$

Proof (i) Let $\frac{1}{p}:=\frac{1}{2}+\frac{\delta}{2}$ and $\frac{1}{q}:=\frac{\delta}{2}$. Then $\frac{1}{p}+\frac{1}{2}=1+\frac{\delta}{2}=1+\frac{1}{q}$ and hence by Young's inequality

$$
\left\|\frac{a}{\langle k\rangle^{\delta}} * b\right\|_{\ell^{q}} \leq C\left\|\frac{a}{\langle k\rangle^{\delta}}\right\|_{\ell^{p}}\|b\|_{\ell^{2}} .
$$

As $p=\frac{2}{1+\delta}>\frac{2}{1+2 \delta}$, Lemma B. 1 (i) can be applied,

$$
\left\|\frac{a}{\langle k\rangle^{\delta}}\right\|_{\ell^{p}} \leq C\|a\|_{\ell^{2}},
$$

and as $q=\frac{2}{\delta}<\infty$, Lemma B. 1 (ii) gives

$$
\left\|\frac{1}{\langle k\rangle}\left(\frac{a}{\langle k\rangle^{\delta}} * b\right)\right\|_{\ell^{1}} \leq C\|a\|_{\ell^{2}}\|b\|_{\ell^{2}}
$$

as claimed.
(ii) By Lemma B.1, for $\frac{1}{p}:=\frac{1}{2}+\frac{2 \delta}{3}\left(<\frac{1}{2}+\delta\right)$

$$
\left\|\frac{a}{\langle k\rangle^{\delta}}\right\|_{\ell^{p}} \leq C\|a\|_{\ell^{2}} .
$$

By Young's inequality with $2 \cdot \frac{1}{p}=1+\frac{1}{q}$ or $\frac{1}{q}=\frac{4 \delta}{3}<1$ (as $0<\delta \leq \frac{1}{2}$ )

$$
\left\|\frac{a}{\langle k\rangle^{\delta}} * \frac{b}{\langle k\rangle^{\delta}}\right\|_{\ell a} \leq C\|a\|_{\ell^{2}}\|b\|_{\ell^{2}}
$$

By Lemma B. 1 (ii) with $\eta=1-\delta$ (hence $1 \leq q=\frac{3}{4 \delta}<\frac{1}{\delta}=\frac{1}{1-\eta}$ )

$$
\left\|\frac{1}{\langle k\rangle^{1-\delta}}\left(\frac{a}{\langle k\rangle^{\delta}} * \frac{b}{\langle k\rangle^{\delta}}\right)\right\|_{\ell^{1}} \leq C\|a\|_{\ell^{2}}\|b\|_{\ell^{2}}
$$

where $1 \leq C \equiv C_{\delta}<\infty$ depends only on $\delta$.

## References

[AKNS] M.I. Ablowitz, D.I. Kaup, A.C. Newell, H.Segur: The inverse scattering transform - Fourier analysis for nonlinear problems. Stud. Appl. Math. 54 (1974), p. 249-315.
[FT] L.D. Faddeev, L.A. Takhtajan: Hamiltonian Methods in the Theory of solitons. Springer Verlag (1987).
[GG] B. Grébert, J.C. Guillot: Gaps of one dimensional periodic AKNS systems. Forum Math. 5 (1993), p. 459-504.
[Gre] B. Grébert: Problèmes spectraux inverses pour les systèmes AKNS sur la droite réelle. Thèse de l'Université de Paris-Nord, 1990.
[GK] B. Grébert, T. Kappeler: Gap estimates of the spectrum of the Zakhavov-Shabat system. Preprint (1997).
[GKM] B. Grébert, T. Kappeler, B. Mityagin: Gap estimates of the spectrum of the Zakhavov-Shabat system. Appl. Math. Lett. 11 (1998), p. 95-97.
[KM] T. Kappeler, B. Mityagin: Estimates for periodic and Dirichlet eigenvalues of the Schrödinger operator. Preprint.
[LS] B.M. Levitan, I.S. Sargsjan: Operators of Sturm-Liouville and Dirac. Nauka, Moscow, 1988.
[LM] Y. Li, D.W. McLaughlin: Morse and Melnikov functions for NLS PDE's. Com. Math. Phys. 162 (1994), p. 175-214.
[MA] Y.C. Ma, M.I. Ablowitz: The periodic cubic Schrödinger equation. Stud. Appl. Math. 65 (1981), p. 113-158.
[Ma] V.A. Marchenko: Sturm-Liouville operators and applications. Operator theory: Advances and Applications 22, Birkhäuser, 1986.
[Mis] T.V. Misyura: Properties of the spectra of periodic and antiperiodic boundary value problems generated by Dirac operators I, II. (In Russian), Teor. Funktsii Funktsii Anal. i. Prilozhen 30 (1978), p. $90-$ 101; 31 (1979), p. 102 - 109; and Finite-zone Dirac operators. (In Russian), Teor. Funktsii Funktsional Anal. i. Prilozhen 33 (1980), p. 107-111.
[ST] J.J. Sansuc, V. Tkachenko: Spectral properties of non-selfadjoint Hill's operators with smooth potentials. In A. Boutet de Monvel and V. Marčenko (eds.), Algebraic and geometric methods in mathematical physics, p. 371-385, Kluwer, 1996.
[Ta1] V. Tkachenko: Non-selfadjoint periodic Dirac operators. Preprint.
[Ta2] V. Tkachenko: Non-selfadjoint periodic Dirac operatos with finite band spectrum. Preprint.
[ZS] V. Zakharov, A. Shabat: A scheme for integrating nonlinear equations of mathematical physics by the method of the inverse scattering problem. Functional Anal. Appl. 8 (1974), p. 226-235.

