

SEMICLASSICAL BEHAVIOUR OF QUANTUM EIGENSTATES ON \mathbb{T}^n

We let $n \geq 1$ and we consider the n -dimensional torus $\mathbb{T}^n = \mathbb{R}^n/2\pi\mathbb{Z}^n$. Given $k \in \mathbb{Z}^n$, we set

$$e_k(x) := \frac{e^{ik \cdot x}}{(2\pi)^{\frac{n}{2}}}.$$

Given $u \in \mathcal{D}'(\mathbb{T}^n)$, we write its Fourier decomposition:

$$u(x) = \sum_{k \in \mathbb{Z}^n} \widehat{u}(k) e_k(x),$$

with, for every $k \in \mathbb{Z}^n$, $\widehat{u}(k) := \langle u, e_{-k} \rangle_{\mathcal{D}' \times \mathcal{C}^\infty}$. Given $a(x, \xi)$ in $\mathcal{C}^\infty(\mathbb{T}^n \times \mathbb{R}^n)$, we also set

$$a(x, \xi) = \sum_{k \in \mathbb{Z}^n} \widehat{a}(k, \xi) e_k(x),$$

where

$$\widehat{a}(k, \xi) := \int_{\mathbb{T}^n} a(x, \xi) e_{-k}(x) dx.$$

For $\hbar > 0$, we define the Weyl quantization of a by its action on the basis $(e_k)_{k \in \mathbb{Z}^n}$:

$$\text{Op}_\hbar^w(a) e_k(x) = \frac{1}{(2\pi)^{\frac{n}{2}}} \sum_{l \in \mathbb{Z}^n} e_l(x) \widehat{a} \left(l - k, \frac{\hbar(k+l)}{2} \right).$$

(1) (a) Show that, if $a \in \mathcal{C}^\infty(\mathbb{T}^n)$, then, for every u in $\mathcal{C}^\infty(\mathbb{T}^n)$,

$$\text{Op}_\hbar^w(a) u = a \times u.$$

(b) Show that, if $a(x, \xi) = \xi_1$, then, for every u in $\mathcal{C}^\infty(\mathbb{T}^n)$,

$$\text{Op}_\hbar^w(a) u = -i\hbar \partial_{x_1} u.$$

(c) Suppose that, for every $\alpha \in \mathbb{Z}^n$, $\|\partial^\alpha a\|_\infty < +\infty$. Using Plancherel equality, show that

$$\text{Op}_\hbar^w(a) : L^2(\mathbb{T}^n) \rightarrow L^2(\mathbb{T}^n)$$

defines a bounded operator.

(d) Let $(k_m)_{m \geq 1}$ be a sequence in \mathbb{Z}^n such that $\hbar_m = \|k_m\|^{-1} \rightarrow 0$ and that $\frac{k_m}{\|k_m\|} \rightarrow \theta$. Determine the accumulation point (in $\mathcal{D}'(\mathbb{T}^n \times \mathbb{R}^n)$) of the sequence

$$w_{\hbar_m} : a \in \mathcal{C}_c^\infty(\mathbb{T}^n \times \mathbb{R}^n) \mapsto \langle e_{k_m}, \text{Op}_{\hbar_m}^w(a) e_{k_m} \rangle_{L^2}, \quad m \geq 1.$$

- (2) (a) Show that
- $\psi_\lambda \in \mathcal{D}'(\mathbb{T}^n)$
- solves

$$-\Delta\psi_\lambda = \lambda^2\psi_\lambda, \quad \|\psi_\lambda\|_{L^2} = 1,$$

if and only if

$$\psi_\lambda = \sum_{\|k\|=\lambda} \widehat{\psi}_\lambda(k)e_k, \quad \sum_{\|k\|=\lambda} |\widehat{\psi}_\lambda(k)|^2 = 1.$$

- (b) Show that, as
- $\lambda \rightarrow +\infty$
- ,

$$N(\lambda, 2\lambda) := |\{k \in \mathbb{Z}^n : \lambda < \|k\| \leq 2\lambda\}| \sim \text{Vol}(B_n(0, 1))(2^n - 1)\lambda^n,$$

where $B_n(0, 1) := \{x \in \mathbb{R}^n : \sum_j x_j^2 \leq 1\}$. *This is the Weyl law.*

- (3) Let
- $a \in \mathcal{C}^\infty(\mathbb{T}^n \times \mathbb{R}^n)$
- all of whose derivatives are bounded.

- (a) Using Plancherel equality, show that

$$\left\| \text{Op}_{\|k\|^{-1}}^w(a)e_k \right\|^2 = \frac{1}{(2\pi)^n} \sum_{l \in \mathbb{Z}^n} \left| \widehat{a} \left(l, \frac{k}{\|k\|} \right) \right|^2 + \mathcal{O}(\|k\|^{-1}).$$

- (b) Show that, as
- $\lambda \rightarrow +\infty$

$$\frac{1}{N(\lambda, 2\lambda)} \sum_{\lambda < \|k\| \leq 2\lambda} \left\| \text{Op}_{\|k\|^{-1}}^w(a)e_k \right\|^2 \longrightarrow \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n \times \mathbb{S}^{n-1}} |a(x, \theta)|^2 dx d\sigma(\theta),$$

where σ is the normalized Lebesgue measure on the unit sphere $\mathbb{S}^{n-1} = \partial B_n(0, 1)$.

- (c) Deduce that, as
- $\hbar \rightarrow 0^+$
- ,

$$\frac{1}{N(\hbar)} \sum_{0 < \hbar\|k\| \leq 1} \left\| \text{Op}_{\|k\|^{-1}}^w(a)e_k \right\|^2 \longrightarrow \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n \times \mathbb{S}^{n-1}} |a(x, \theta)|^2 dx d\sigma(\theta),$$

where

$$N(\hbar) := |\{k \in \mathbb{Z}^n : 0 < \hbar\|k\| \leq 1\}|.$$

This is a microlocal Weyl law.

- (4) Let
- a
- be an element in
- $\mathcal{C}^0(\mathbb{T}^n \times \mathbb{S}^{n-1})$
- .

- (a) Show that, for every
- x
- in
- \mathbb{T}^n
- and for
- σ
- a.e.
- θ
- in
- \mathbb{S}^{n-1}
- , one has

$$\lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T a(x + t\theta, \theta) dt = \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n} a(x, \theta) dx.$$

Indication. Start with the case where a is trigonometric polynomial.

- (b) Deduce that, if
- a
- does not depend on the variable
- θ
- , then

$$\lim_{T \rightarrow +\infty} \int_{\mathbb{T}^n \times \mathbb{S}^{n-1}} \left| \frac{1}{T} \int_0^T a(x + t\theta) dt - \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n} a(x) dx \right|^2 dx d\sigma(\theta) = 0.$$

- (5) Let
- a
- be an element in
- $\mathcal{C}^\infty(\mathbb{T}^n)$
- . Show that, for every
- $t \in \mathbb{R}$
- ,

$$e^{-\frac{it\hbar\Delta}{2}} a e^{\frac{it\hbar\Delta}{2}} = \text{Op}_\hbar^w(a_t),$$

where $a_t(x, \xi) = a(x + t\xi)$.

(6) Let $(\psi_j)_{j \geq 1}$ be an orthonormal basis of $L^2(\mathbb{T}^n)$ such that, for every $j \geq 1$,

$$-\Delta \psi_j = \lambda_j^2 \psi_j,$$

with $0 = \lambda_1 < \lambda_2 \leq \dots \leq \lambda_j \leq \lambda_{j+1} \leq \dots$. For $j > 1$, we set $\hbar_j := \lambda_j^{-1}$. Let a in $\mathcal{C}^\infty(\mathbb{T}^n)$ such that $\int_{\mathbb{T}^n} a dx = 0$.

(a) Show that, for every $T > 0$

$$V_\hbar(a) := \frac{1}{N(\hbar)} \sum_{0 < \hbar \lambda_j \leq 1} |\langle \psi_j, a \psi_j \rangle|^2 \leq \frac{1}{N(\hbar)} \sum_{0 < \hbar \lambda_j \leq 1} \left\| \text{Op}_{\hbar_j}^w(\langle a \rangle_T) \psi_j \right\|^2,$$

where $\langle a \rangle_T(x, \xi) = \frac{1}{T} \int_0^T a(x + t\xi) dt$.

(b) Deduce that, for every $T > 0$,

$$V_\hbar(a) \leq \frac{1}{N(\hbar)} \sum_{0 < \hbar \|k\| \leq 1} \left\| \text{Op}_{\|k\|^{-1}}^w(\langle a \rangle_T) e_k \right\|^2.$$

(c) Conclude that

$$\lim_{\hbar \rightarrow 0} V_\hbar(a) = 0.$$

(7) Let a in $\mathcal{C}^\infty(\mathbb{T}^n)$. Deduce from the previous question that there exists $S \subset \mathbb{Z}_+^*$ such that

$$\lim_{N \rightarrow +\infty} \frac{|\{1 \leq j \leq N : j \in S\}|}{N} = 1,$$

and

$$\lim_{j \rightarrow +\infty, j \in S} \int_{\mathbb{T}^n} a(x) |\psi_j(x)|^2 dx = \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n} a(x) dx.$$

(8) Show that S can be chosen such that

$$\lim_{N \rightarrow +\infty} \frac{|\{1 \leq j \leq N : j \in S\}|}{N} = 1,$$

and, for every a in $\mathcal{C}^0(\mathbb{T}^n)$,

$$\lim_{j \rightarrow +\infty, j \in S} \int_{\mathbb{T}^n} a(x) |\psi_j(x)|^2 dx = \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n} a(x) dx.$$

(9) Is it necessary to extract a subsequence?